

**MODELLING
SOLAR AND STELLAR ACTIVITY
DRIVEN BY
TURBULENT DYNAMO EFFECTS AND
HELICITY**

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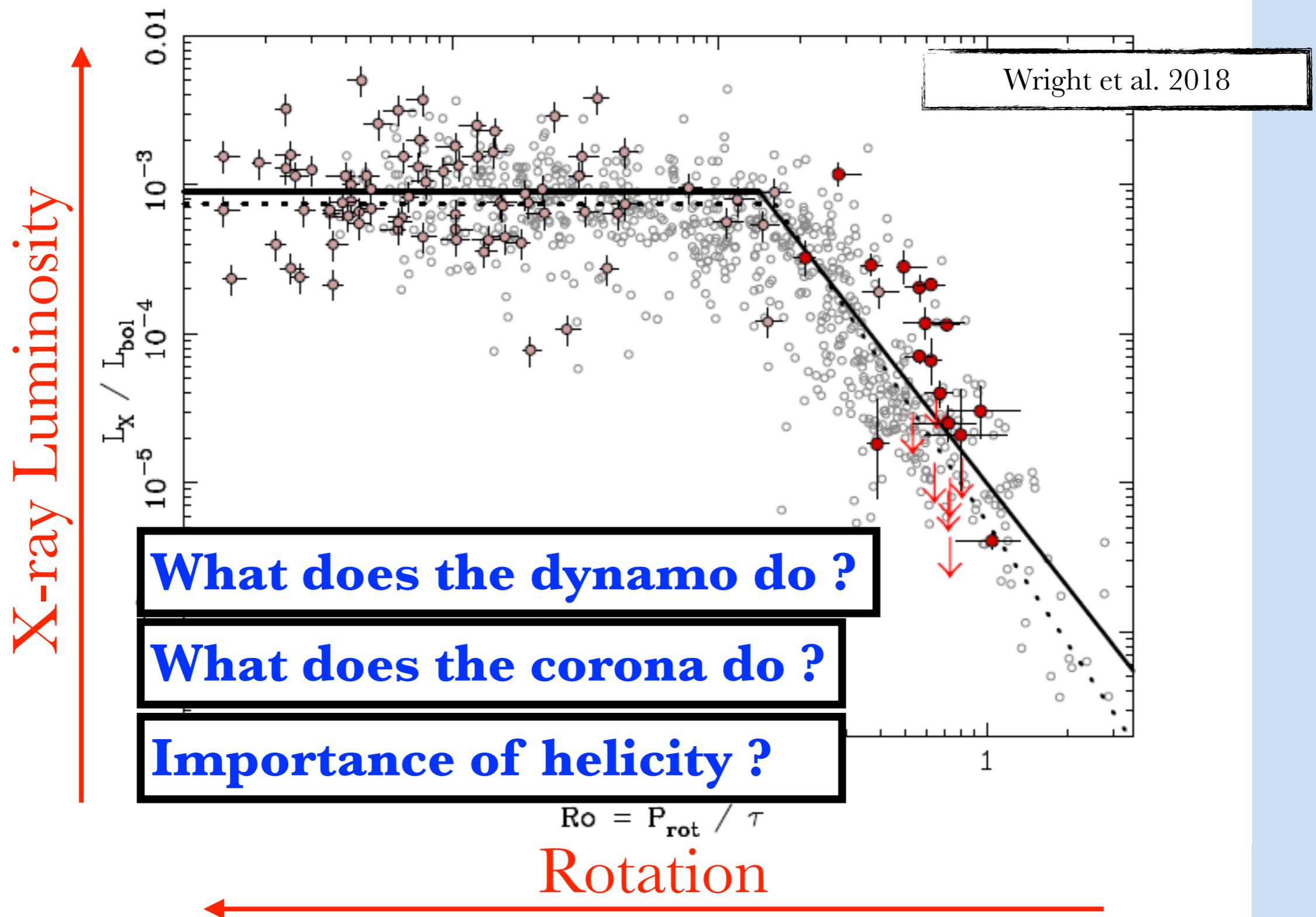
AXEL BRANDENBURG, NORDITA

MATTHIAS RHEINHARDT, AALTO UNIVERSITY

HARDI PETER & JUXHIN ZHULENKU, MPS



Rotation-Activity Relation



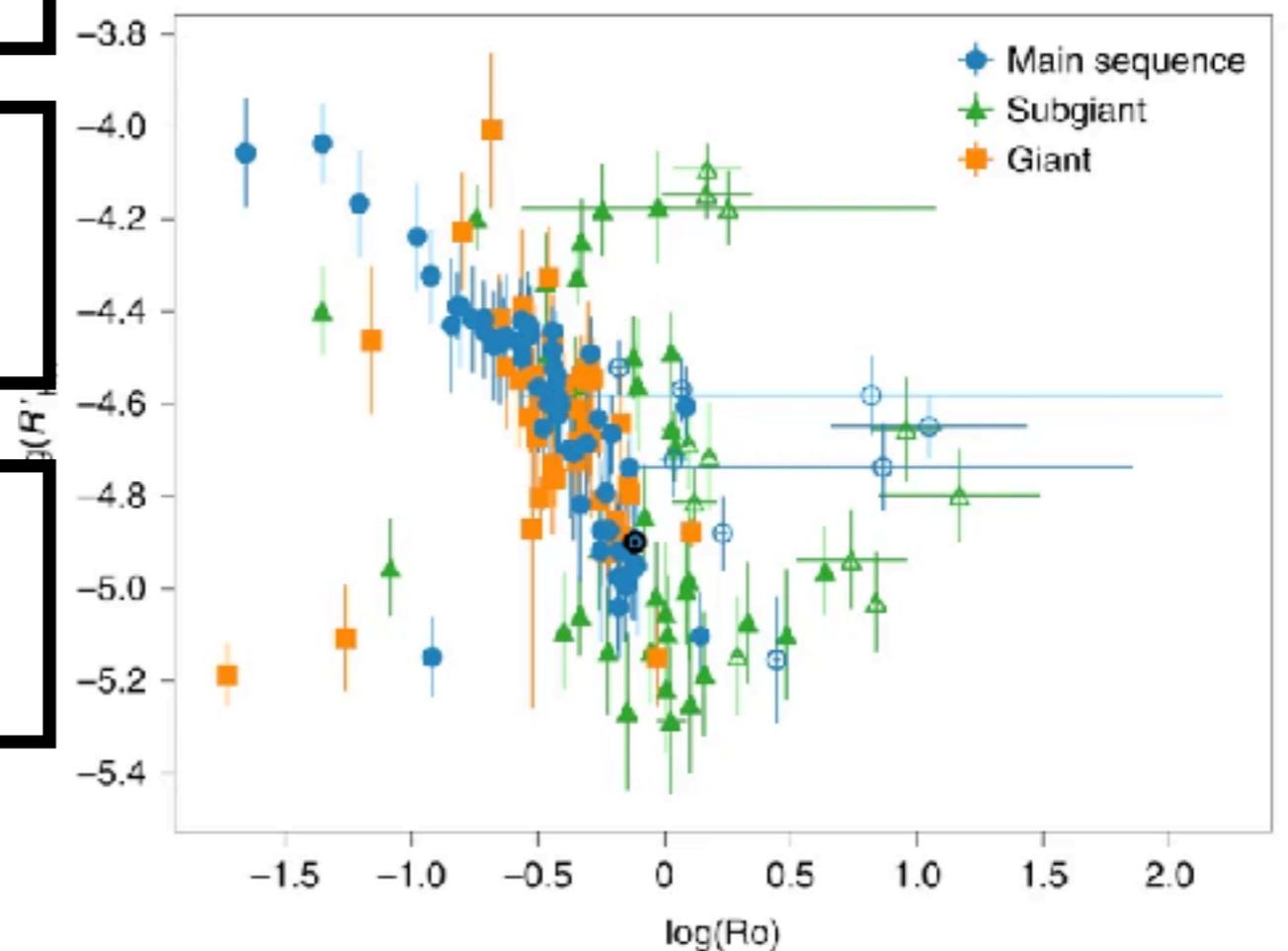
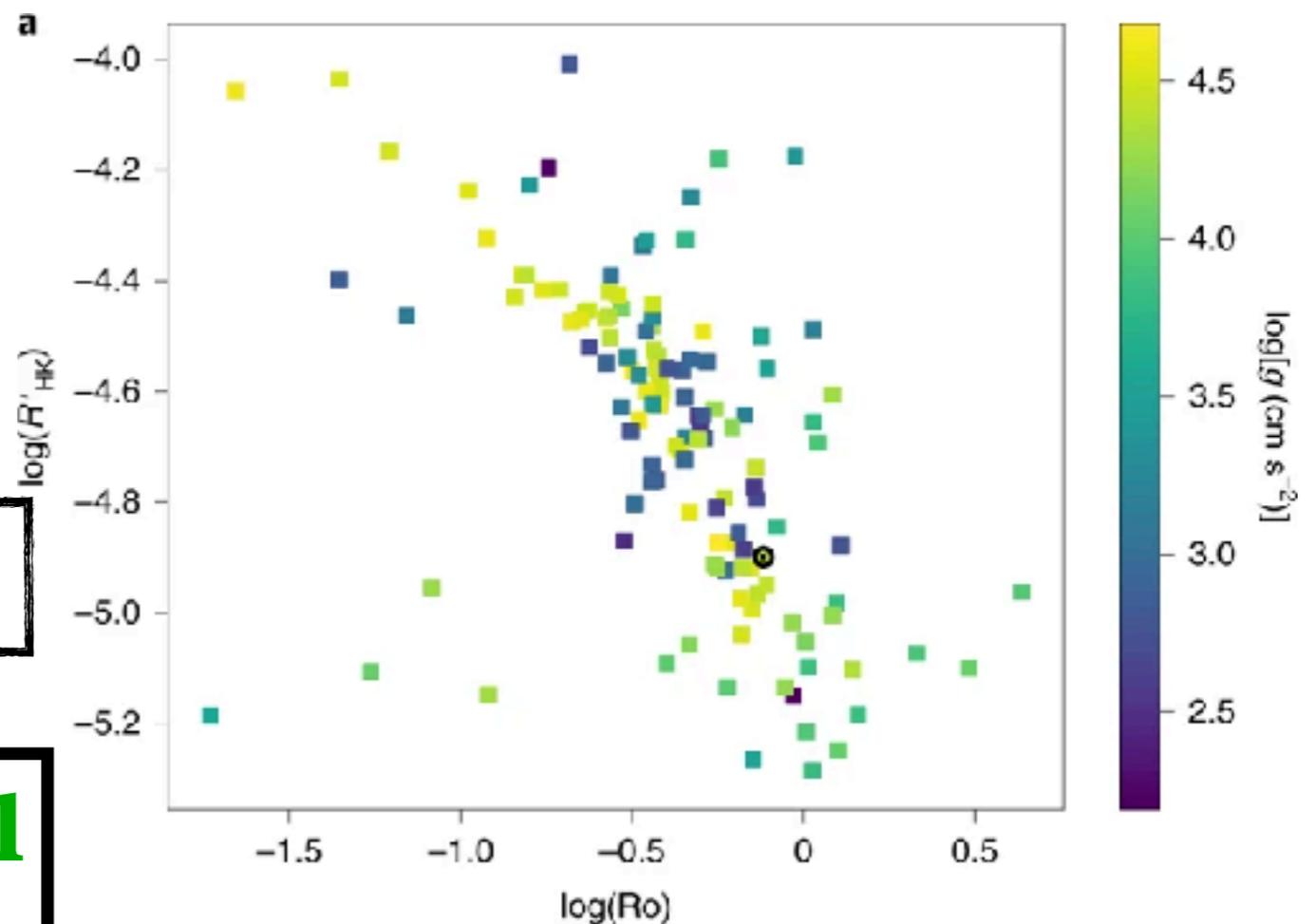
Rotation-Activity Relation II

Lehtinen et al. 2020
Nature Astronomy

MS stars and Giants fall on top of each other

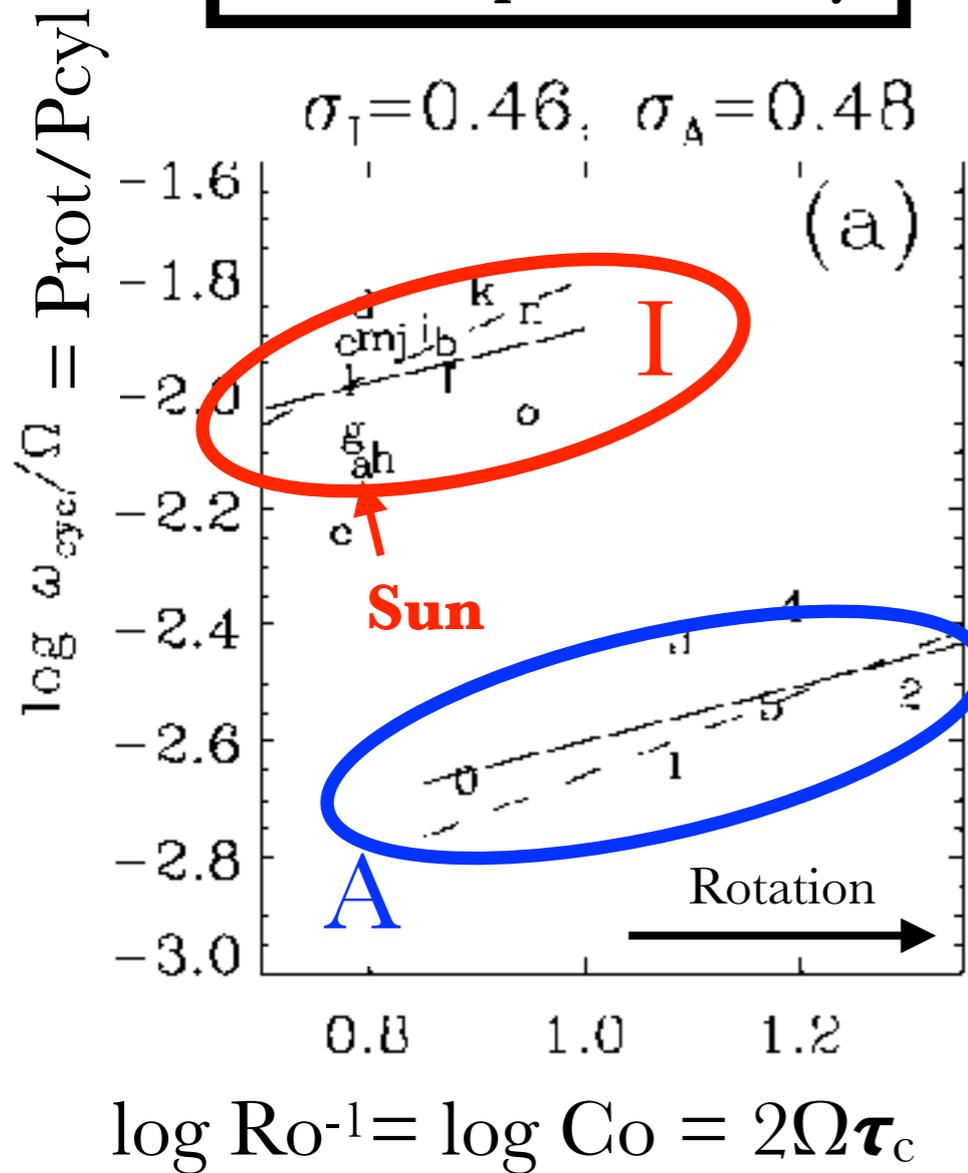
Rossby number the crucial parameter

Convective turbulence important for Dynamo



Activity Cycles

Chromospheric activity



Active branch under debate

(e.g. Distefano et al. 2017, Brandenburg et al. 2017, Reinhold et al. 2017, Boro Saikia et al. 2018)

Inactive branch prevails

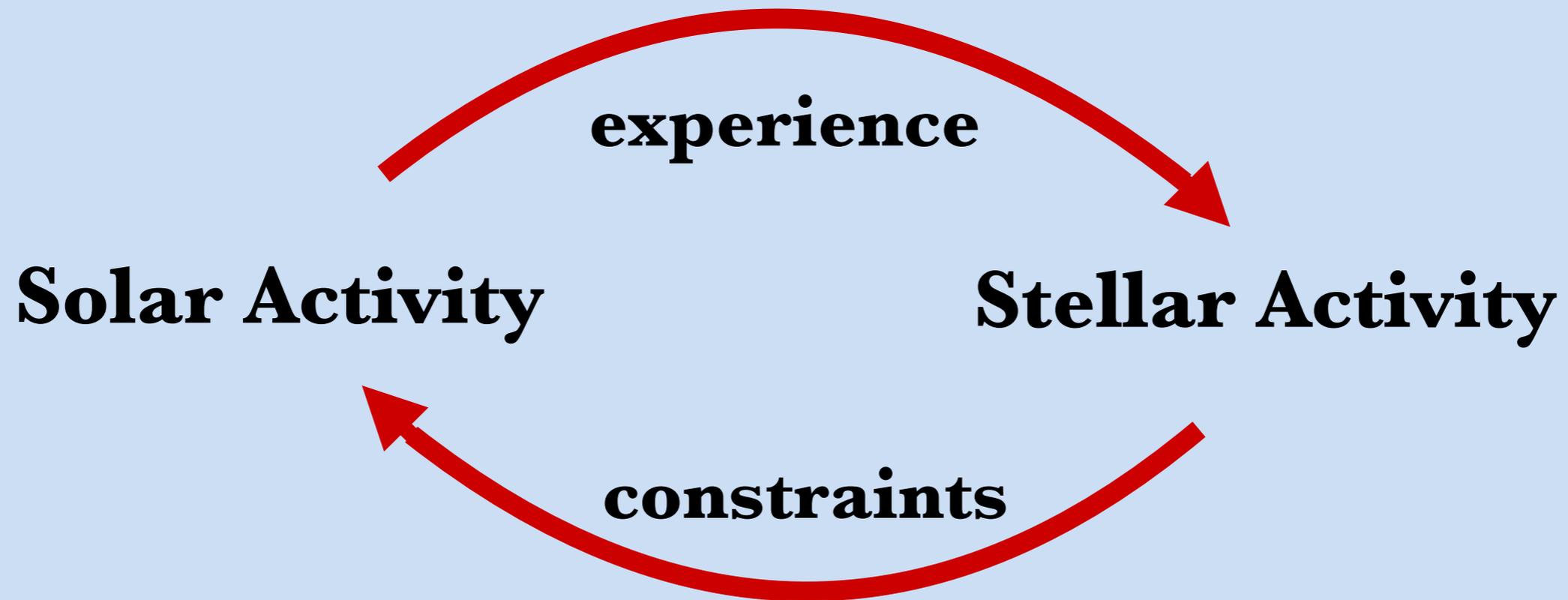
(e.g. Olperets et al. 2018)

$$\text{Prot}/\text{Pcyl} \sim \text{Co}^{0.5} \rightarrow \text{Pcyl} \sim \text{Co}^{-1.5}$$

**Strong constrains
on the solar dynamo**

Brandenburg, Saar & Turpin, 1998

- **Modelling of dynamos, sunspots and the corona**
- **High-resolution and long-term observation**



- **Rotational dependence of cycles & X-rays**
- **Size and distribution of star spots**
- **Dependence on stellar parameters**

Magnetic Helicity



**The glue that connects
dynamoes and coronae
of the Sun and stars**

Helicity in the Sun and Stars

**Alpha effect from
global simulations**

**Magnetic helicity
production**

Magn. helicity fluxes

**Dynamo - Corona
connection**

**Importance of
magnetic helicity
for coronal heating
and X-ray emission**

**Nonalignment of
rotation and gravity**

Kinetic helicity

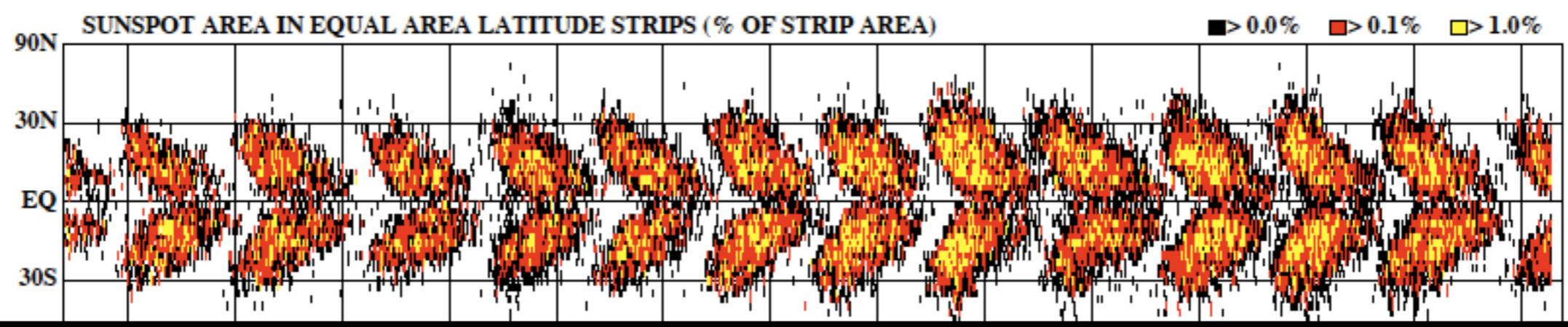
Alpha-effect

**Magnetic helicity
+ catastroph. quenching**

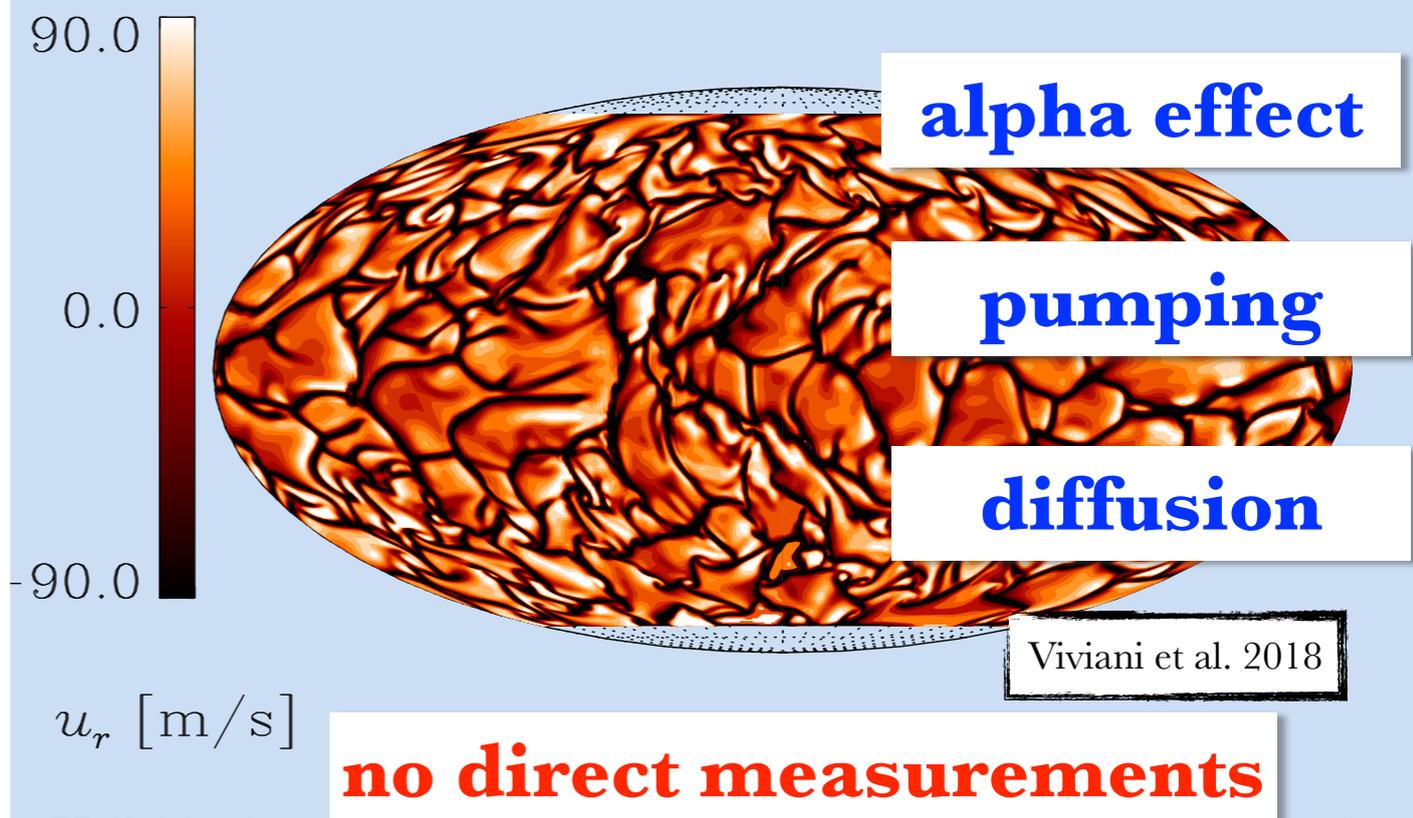
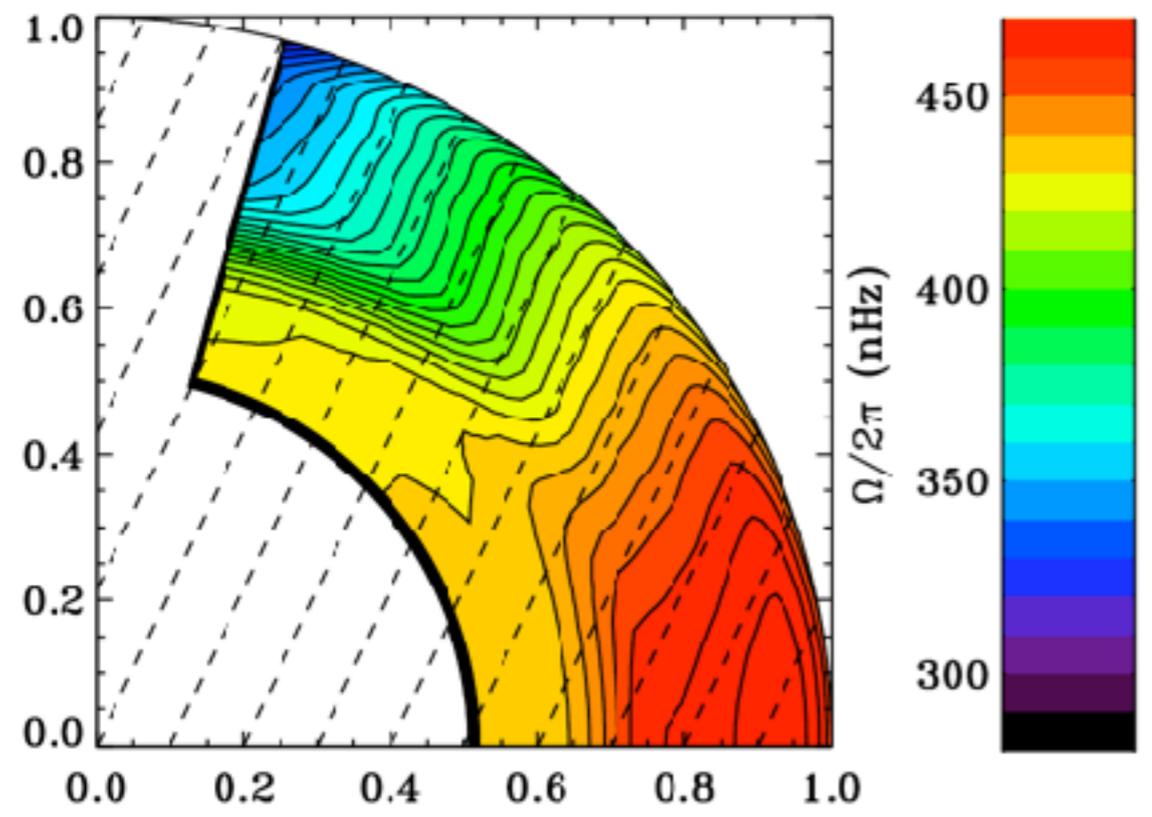
**Space weather
coronal heating**

Solar and Stellar Dynamos

DAILY SUNSPOT AREA AVERAGED OVER INDIVIDUAL SOLAR ROTATIONS



We can use numerical simulations to determine turbulent effects.



Dynamo theory

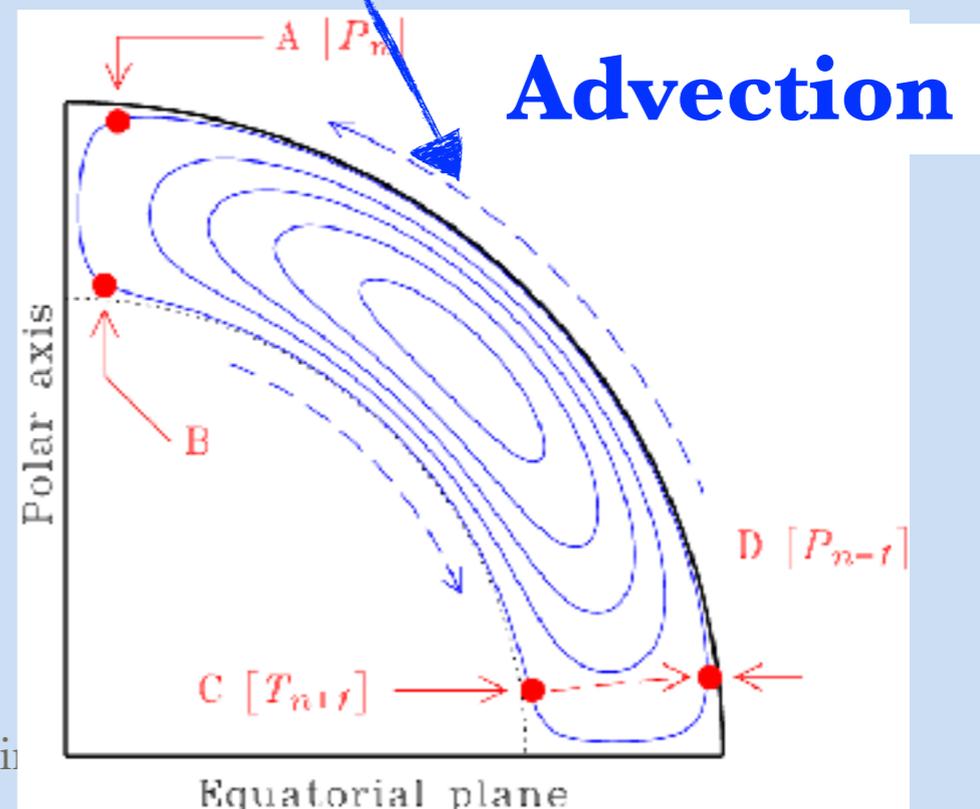
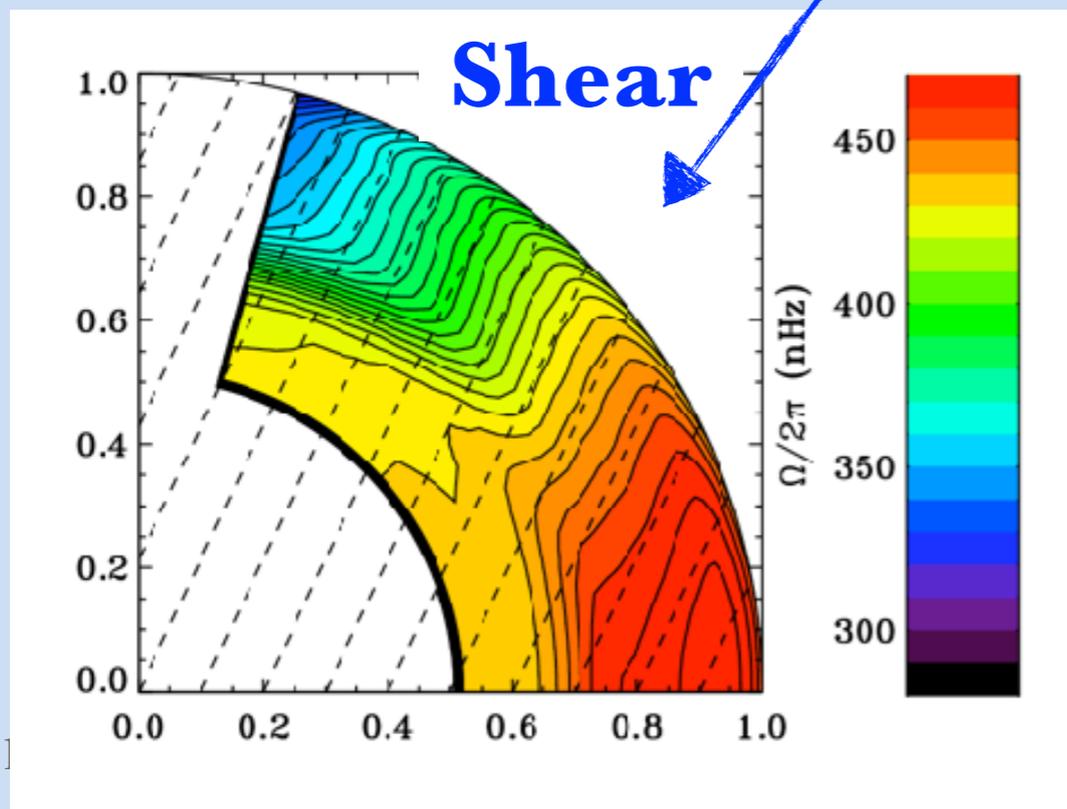
$$\frac{\partial B}{\partial t} = \nabla \times (u \times B) - \nabla \times \eta J$$

$$B = \bar{B} + b'$$

**Electromotive force:
Turbulent dynamo effects**

$$\frac{\partial \bar{B}}{\partial t} = \nabla \times (\bar{U} \times \bar{B} + \overline{u' \times b'}) - \nabla \times \eta \bar{J}$$

$$\nabla \times (\bar{U} \times \bar{B}) = (\bar{B} \cdot \nabla) \bar{U} - \bar{B} (\nabla \cdot \bar{U}) - (\bar{U} \cdot \nabla) \bar{B}$$



Electromotive force I

$$\mathcal{E} = a \cdot \bar{\mathbf{B}} + b \cdot \nabla \bar{\mathbf{B}} + \dots$$

$$\mathcal{E}_i = a_{ij} \bar{B}_j + b_{ijk} \partial_j \bar{B}_k + \dots$$

$$\mathcal{E} = \alpha \cdot \bar{\mathbf{B}} + \gamma \times \bar{\mathbf{B}} - \beta \cdot (\nabla \times \bar{\mathbf{B}}) - \delta \times (\nabla \times \bar{\mathbf{B}}) - \kappa \cdot (\nabla \bar{\mathbf{B}})^{(S)}$$

Alpha effect

Turbulent diffusion

Turbulent pumping

Rädler effect

Electromotive force II

Contributions:

Alpha effect: amplification

„Easy“ approach
scalar, only acts on toroidal field

Turbulent pumping: transport

radial, tuned to fit the Sun

Turbulent diffusion: diffusion

scalar, tuned to fit the Sun

Rädler effect: delta effect

Not considered

+ **Additional turbulent**

Not considered

Schrinner et al. 2005, 2007, 2012

Test-field method: Measuring turbulent dynamo effects

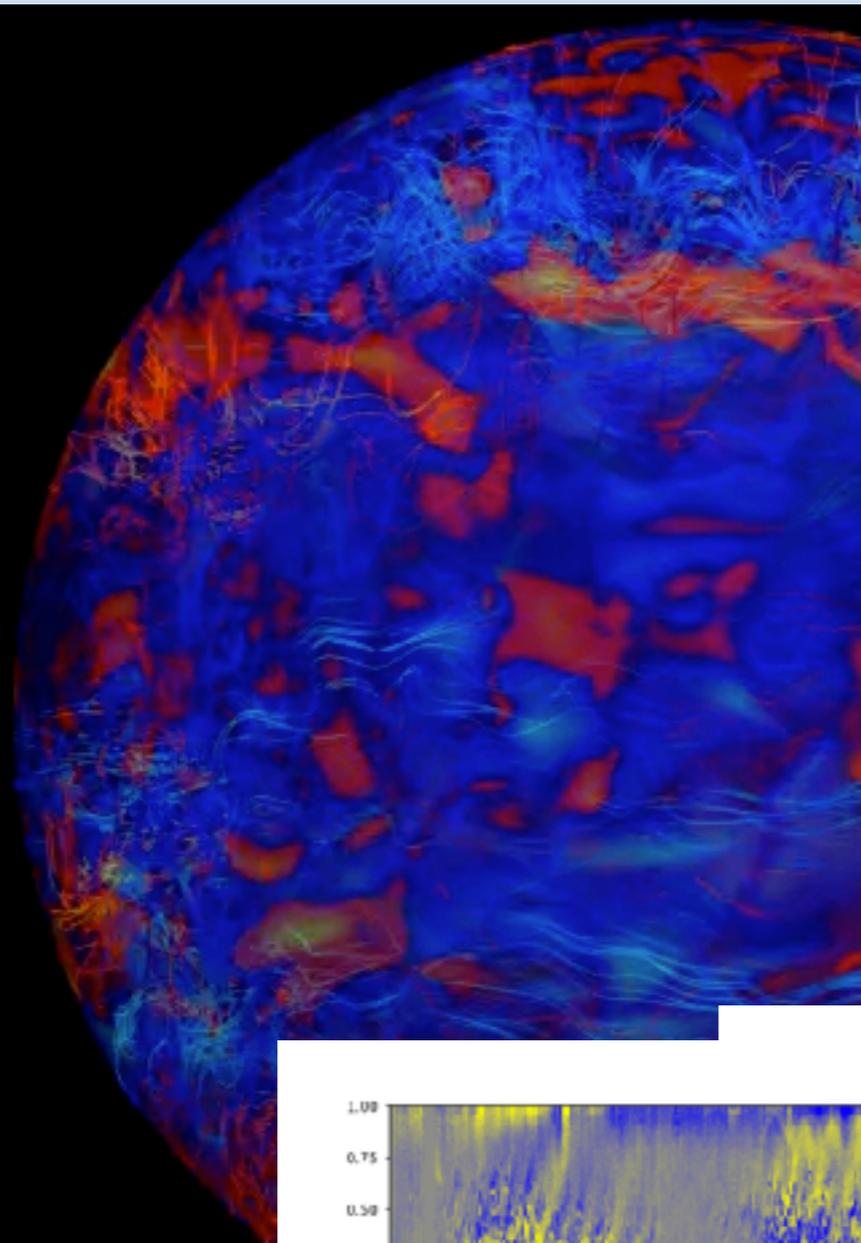
Applying 9 linear independent test-fields, no back reaction

Calculating corresponding electromotive force

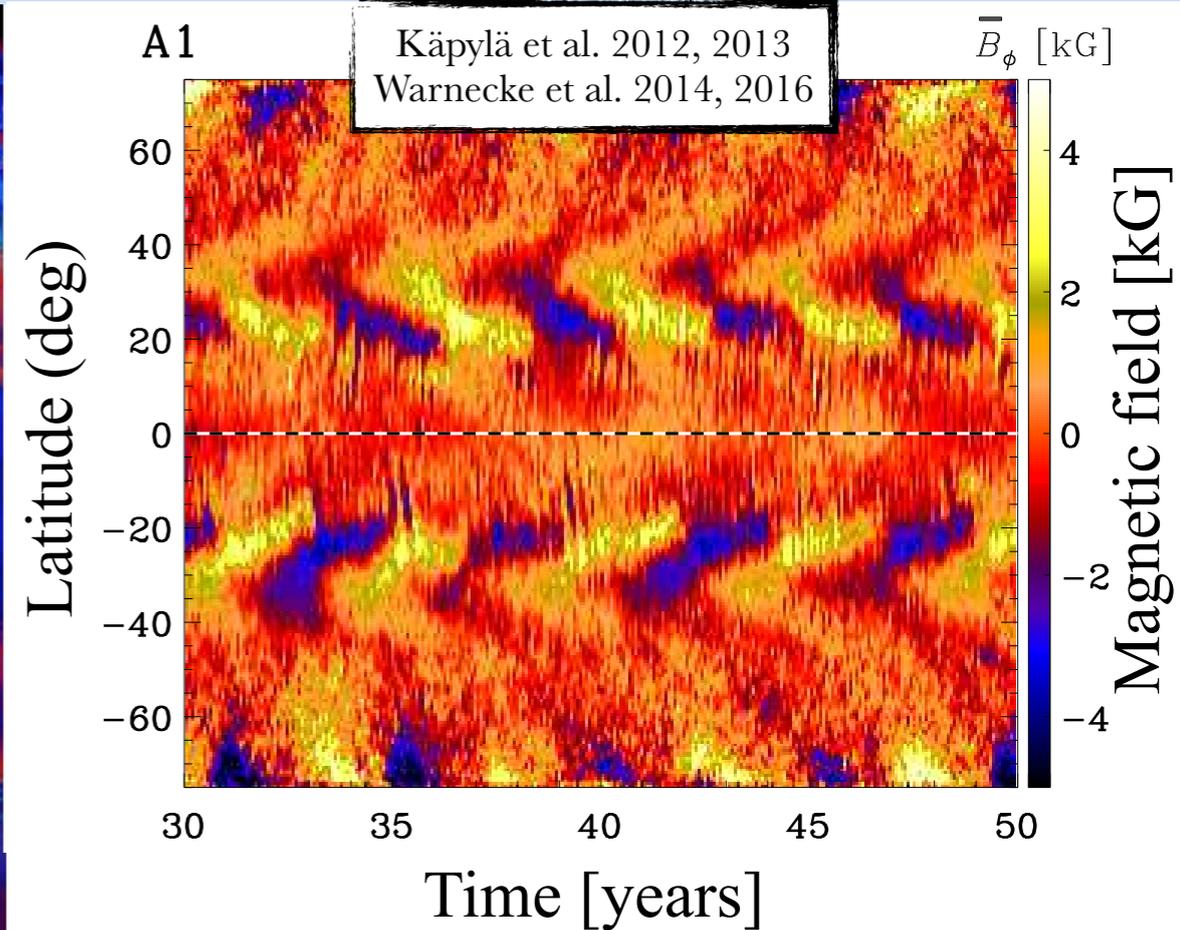
Inverting for turbulent dynamo effects

Global stellar dynamo simulations

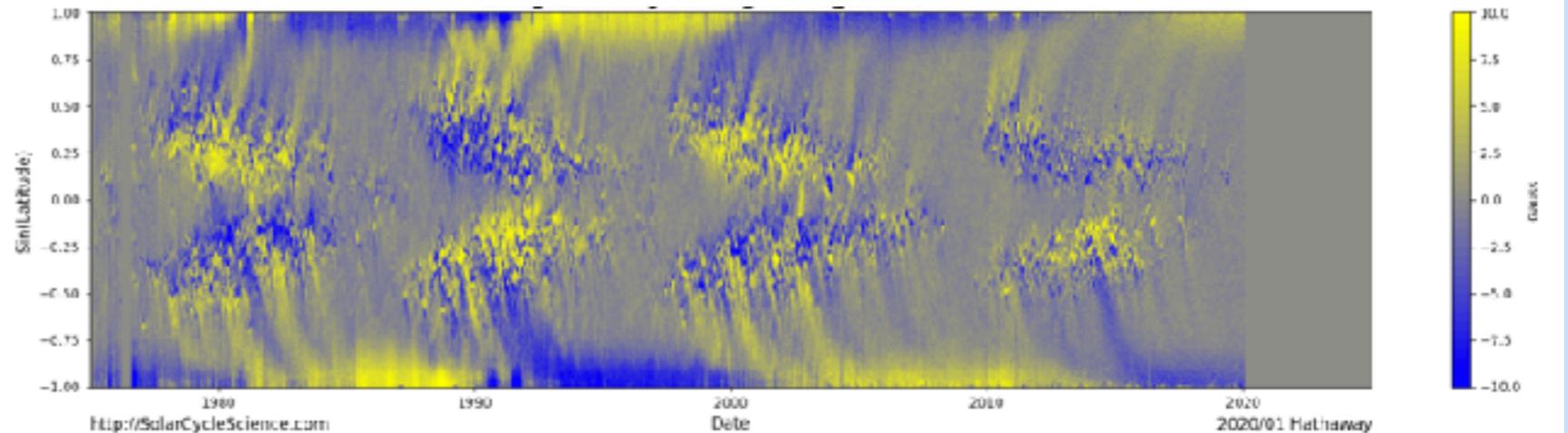
Pencil Code



Viviani et al. 2018

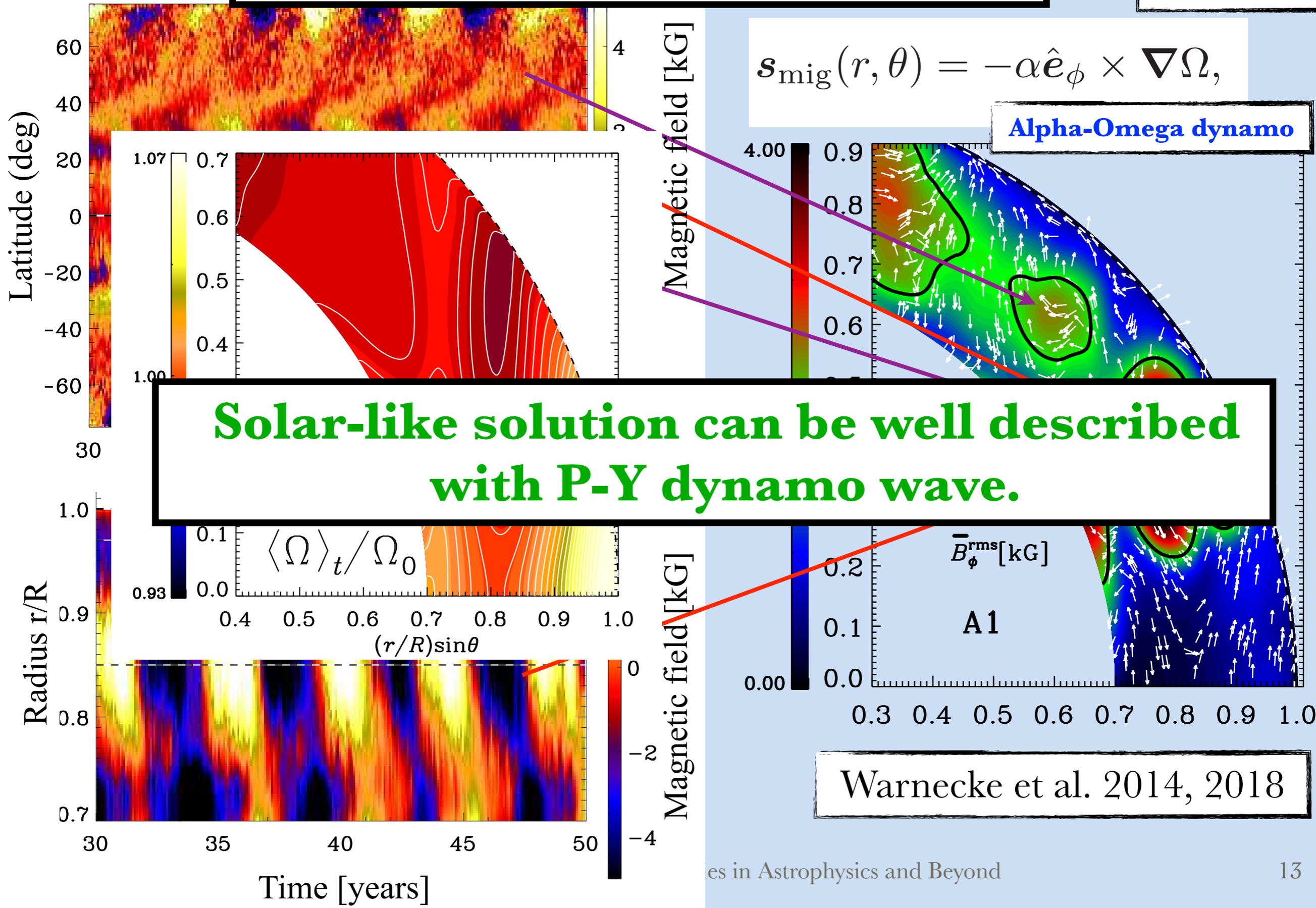


Solar magnetic field



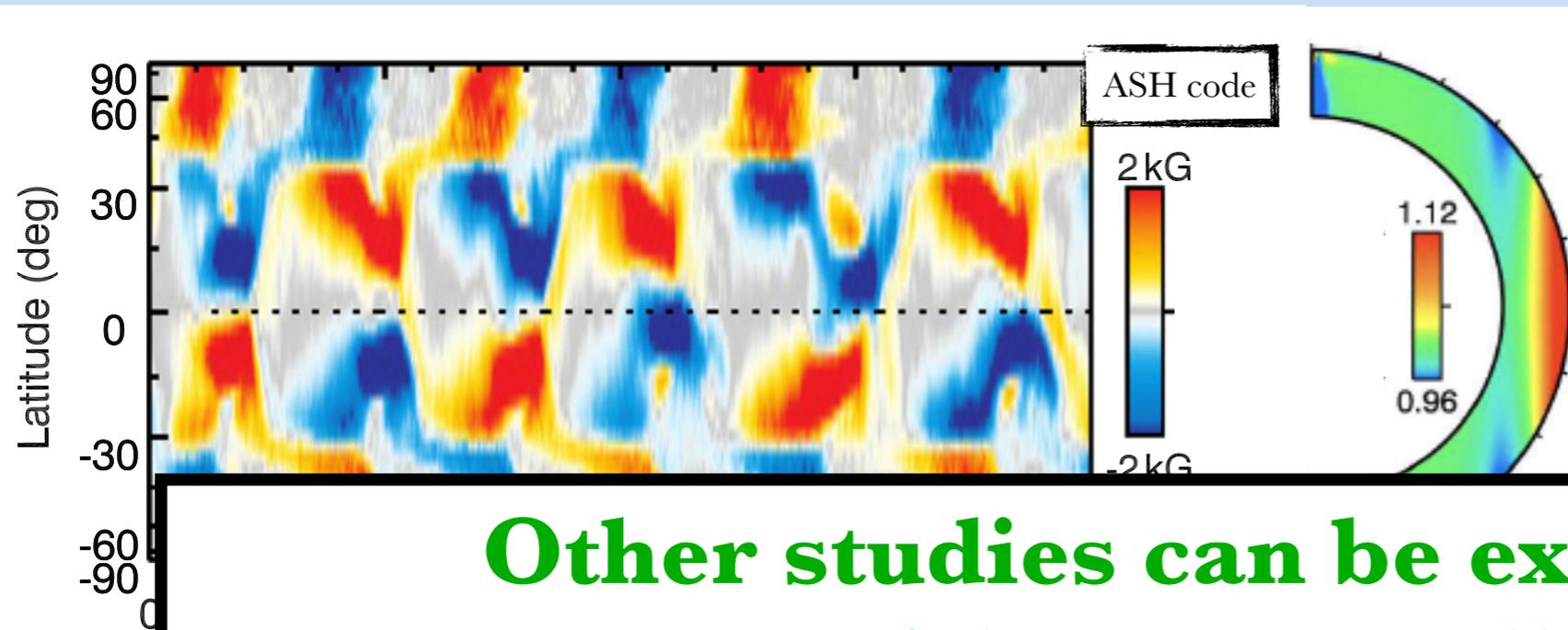
Parker—Yoshimura Rule

Parker 1955
Yoshimura 1975



Warnecke et al. 2014, 2018

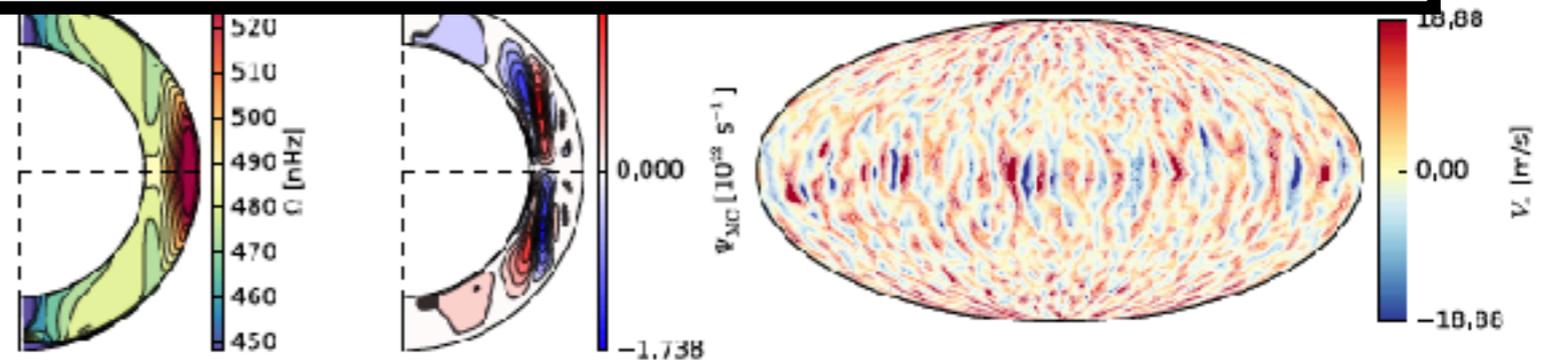
Dynamo cycles of other authors



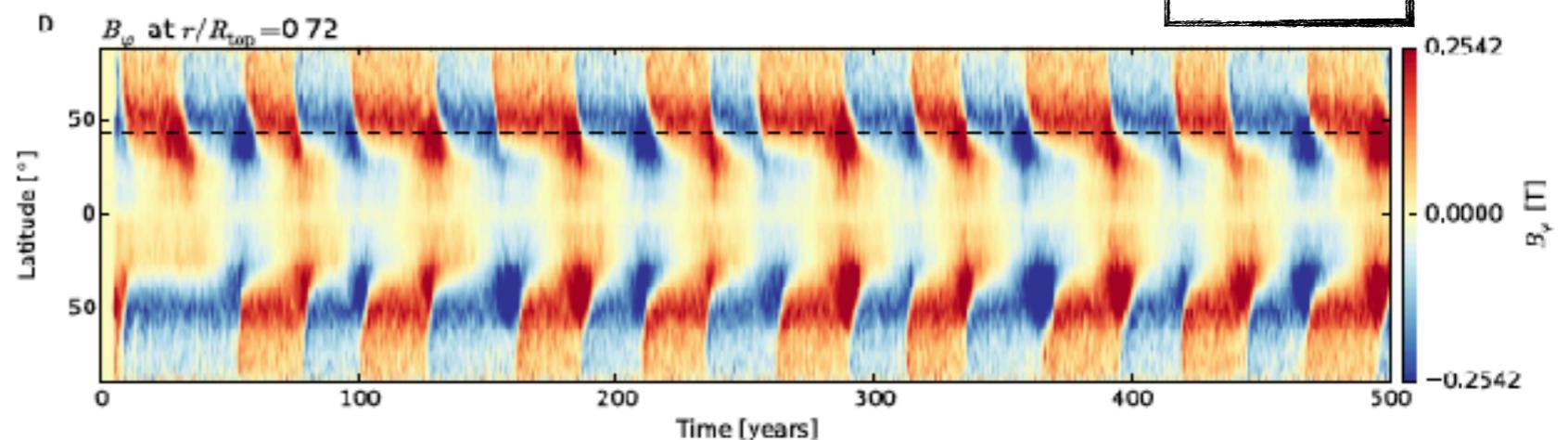
Other studies can be explained with P-Y as well.

2017

Augustso



EULAC code

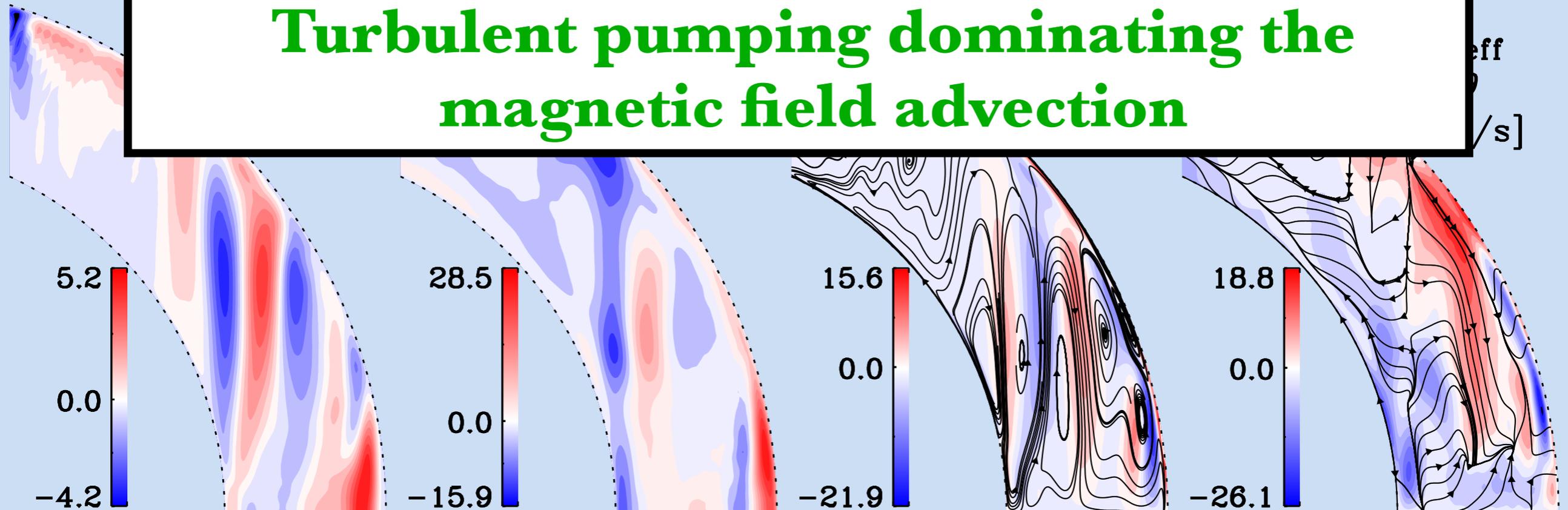


Turbulent pumping

Mean magnetic field only „sees“ the sum of flow and pumping

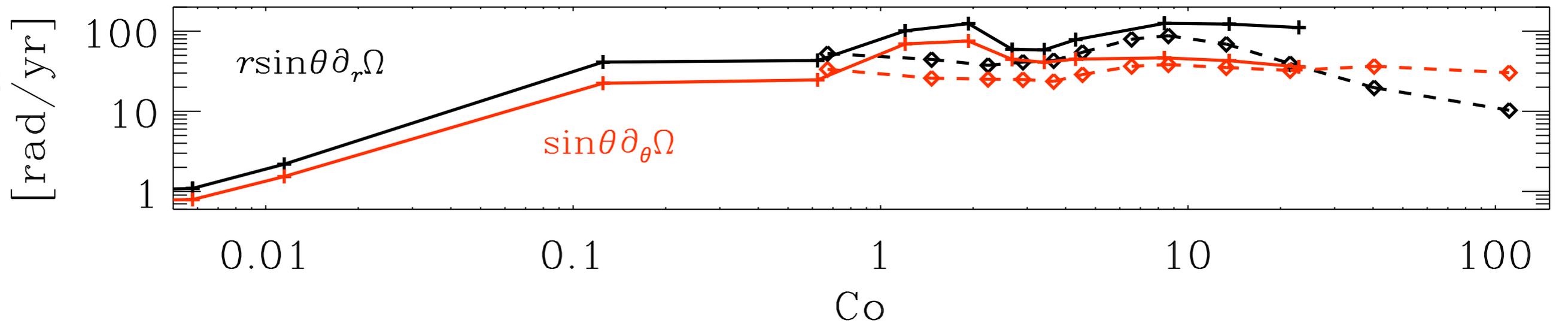
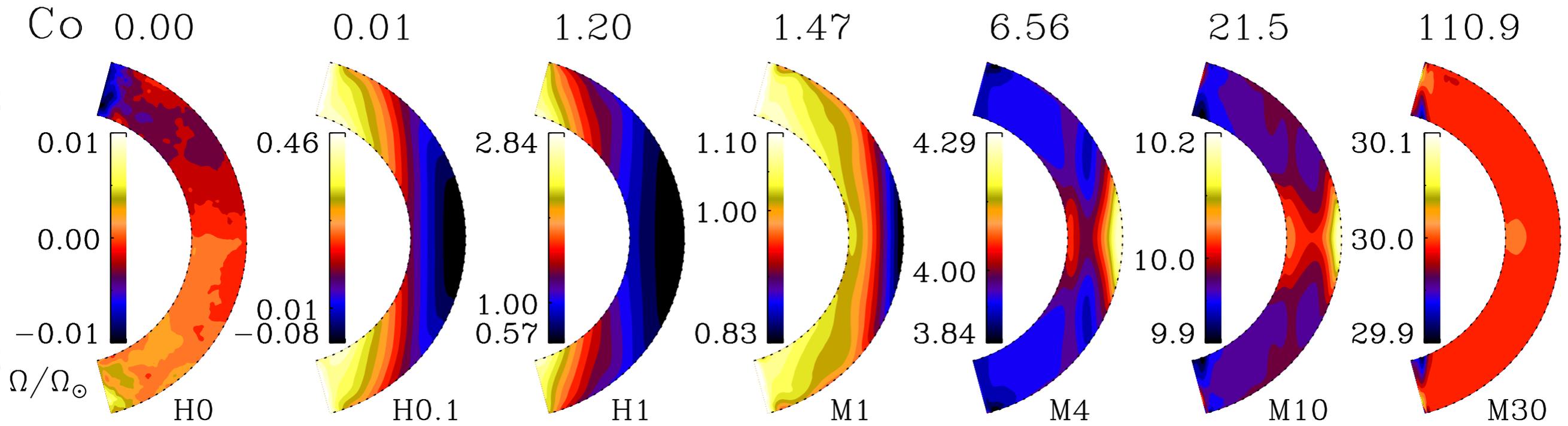
$$\overline{U}_{\text{eff}} = \overline{U} + \gamma$$

Turbulent pumping dominating the magnetic field advection



Warnecke et al. 2018

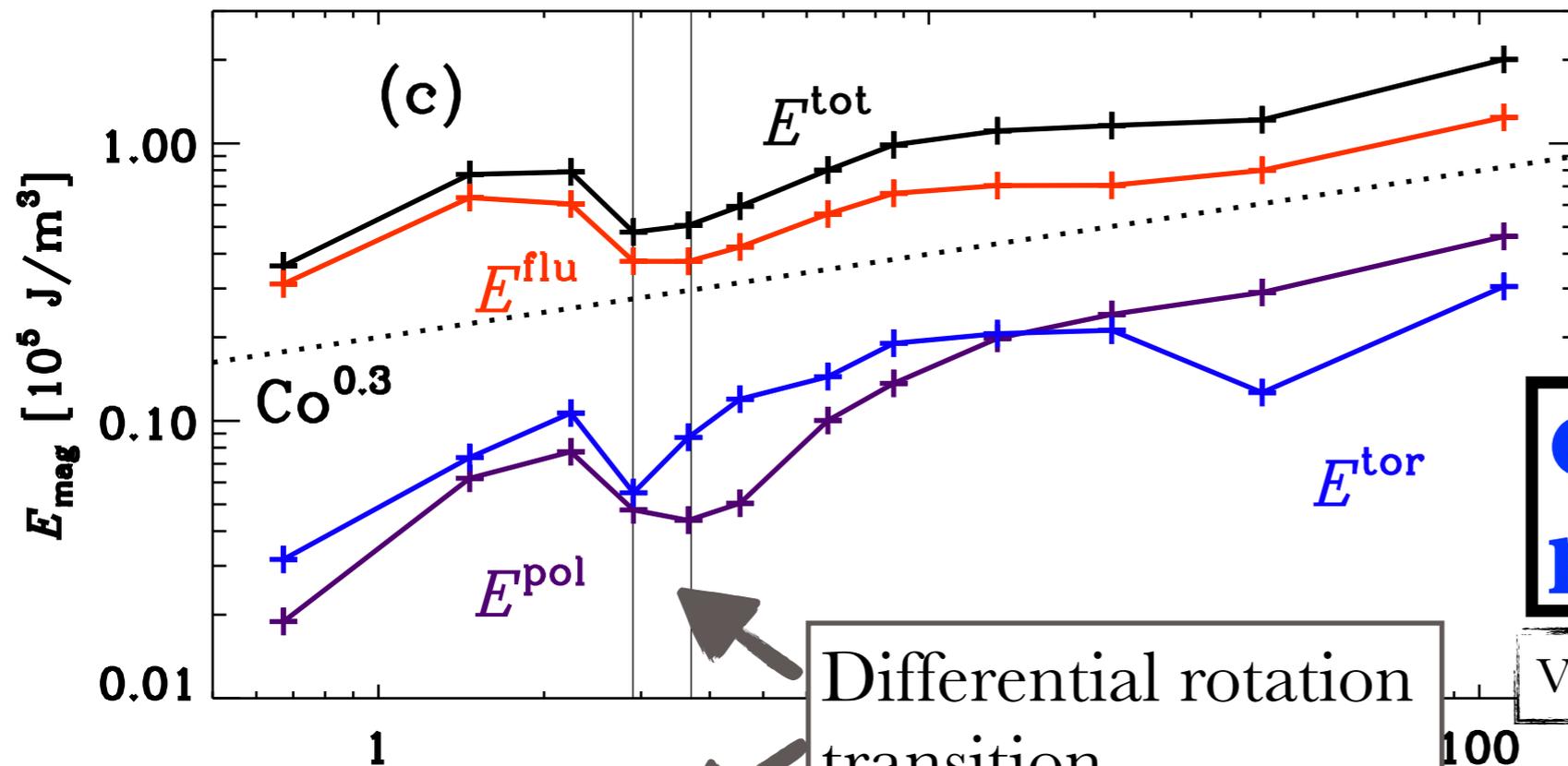
Dynamo simulations - increasing rotation



Data on ZENODO:
DOI: 10.5281/zenodo.3629665

Warnecke and Käpylä, 2020

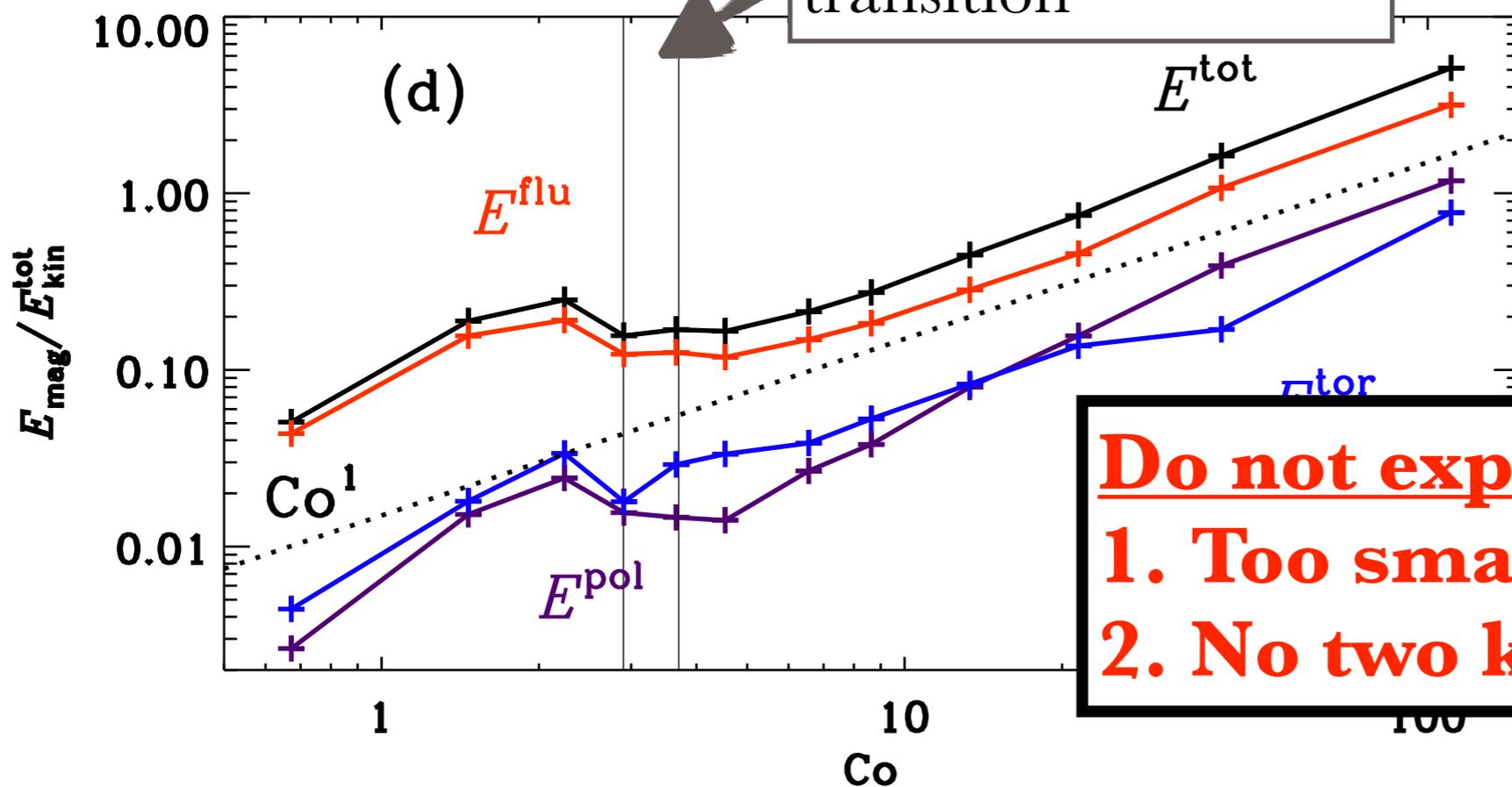
Magnetic energy



Consistent with previous simulations

Viviani et al. 2018

Augustson et al. 2019

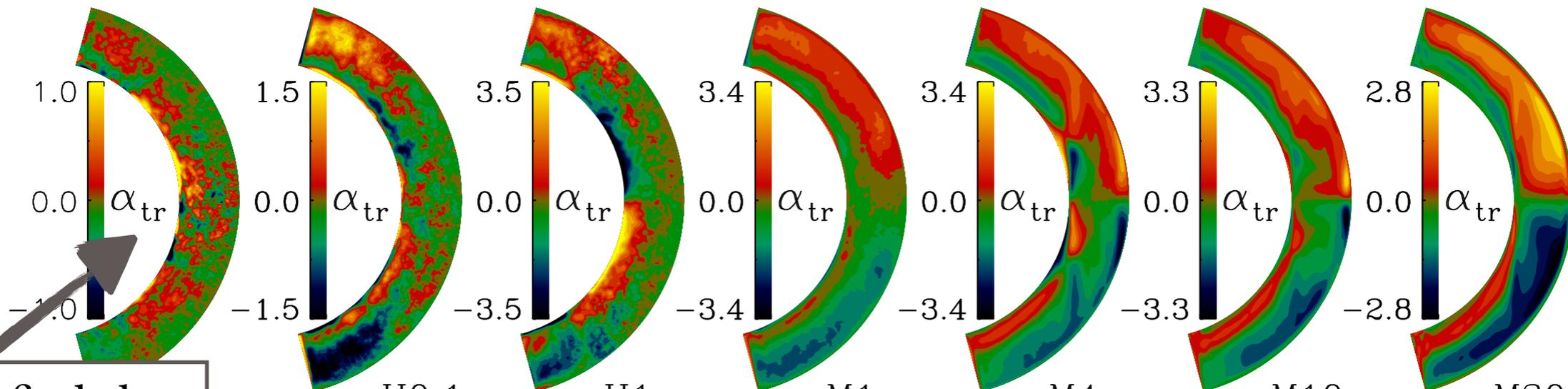


Do not explain observations:

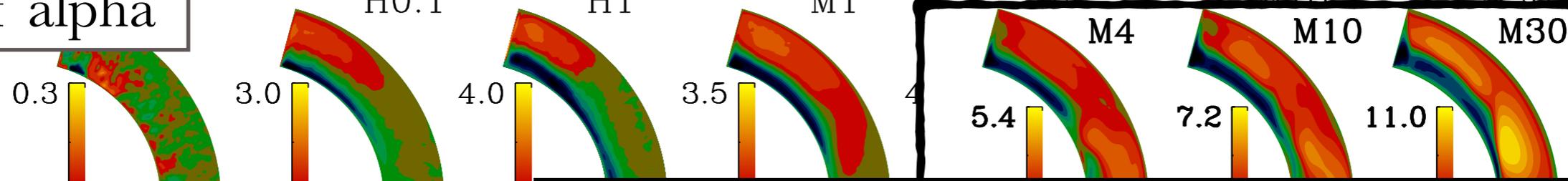
1. Too small slope
2. No two kinds of dependency

Turbulent transport coefficients: alpha

Co 0.00 0.12 1.20 1.46 6.56 21.5 110.9



Trace of alpha



$$\tau_c \sim Co^{-0.5} ?$$

$Co^{-0.5}$:

Viviani et al. 2018

Featherstone & Hindman, 2016

$Co^{-0.3}$:

Chandrasekhar, 1961

$$\alpha_K = -\frac{\tau_c}{3} \overline{\omega' \cdot u'}, \quad \alpha_M = \frac{\tau_c}{3} \overline{j' \cdot b' / \bar{\rho}}, \quad \alpha_{KM} = \alpha_K - \alpha_M$$

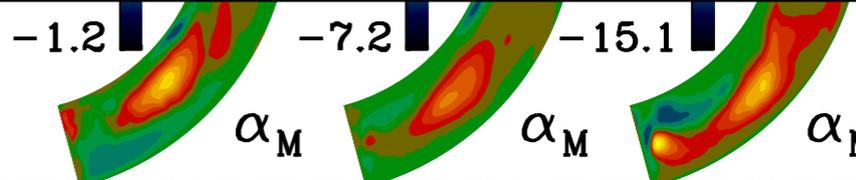
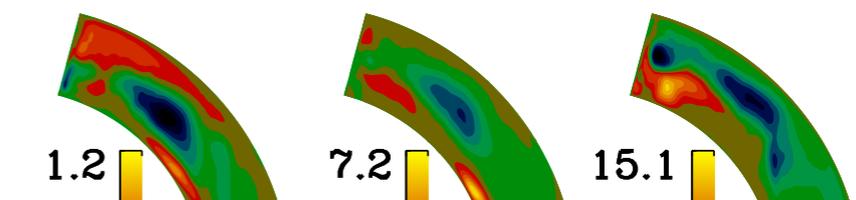
α_{rms} / α_0

Magnetic alpha not important ?

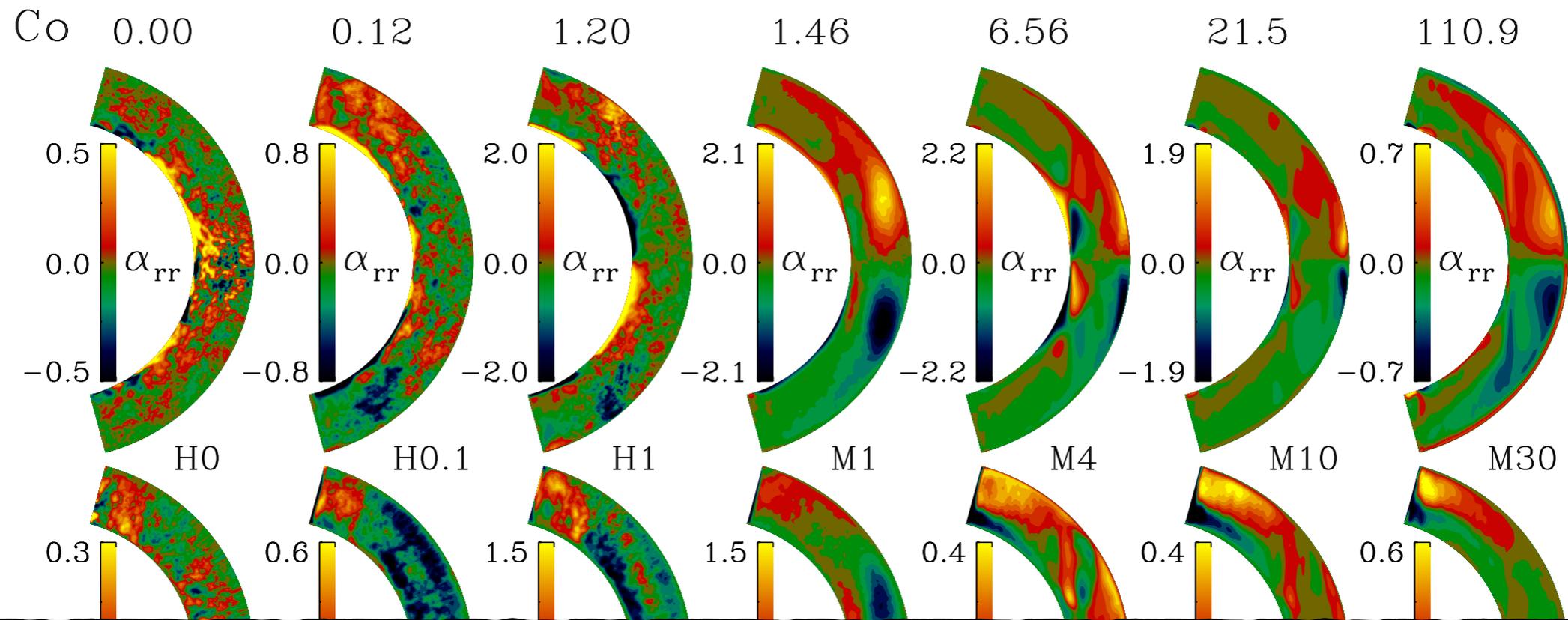
$$\alpha_0 = u'_{rms} / 3.$$

0.01 0.1 1

Co



Turbulent transport coefficients: alpha

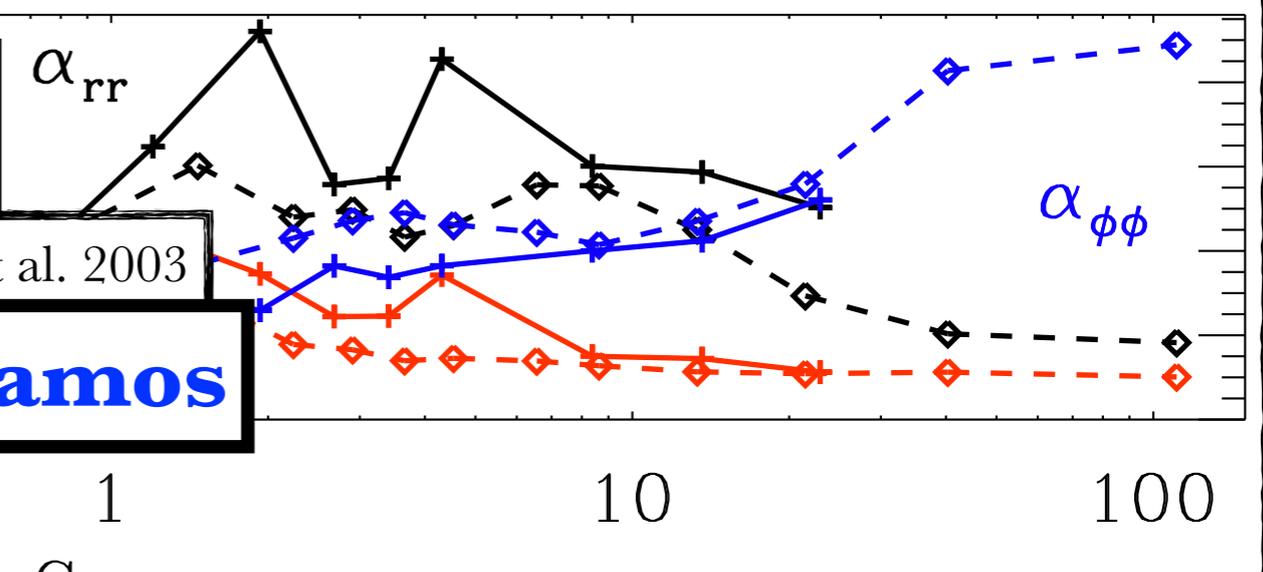


α becomes highly anisotropic

Rädler et al. 1990, Rüdiger & Kitchatinov 1993 & Rüdiger et al. 2003

Explain non-axisymmetric dynamos

Rädler et al. 1990, Pipin 2017



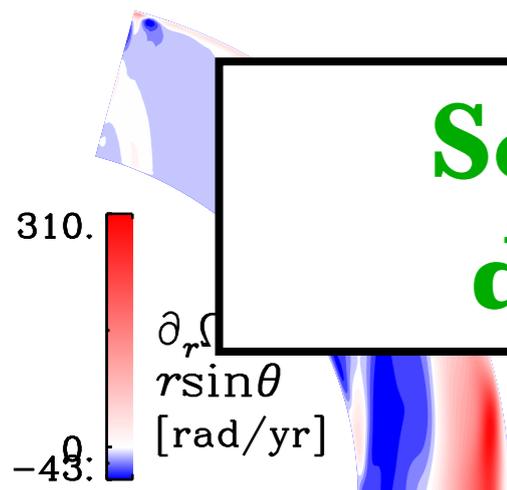
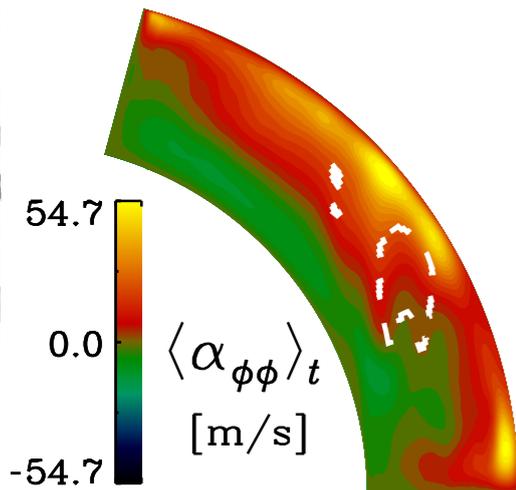
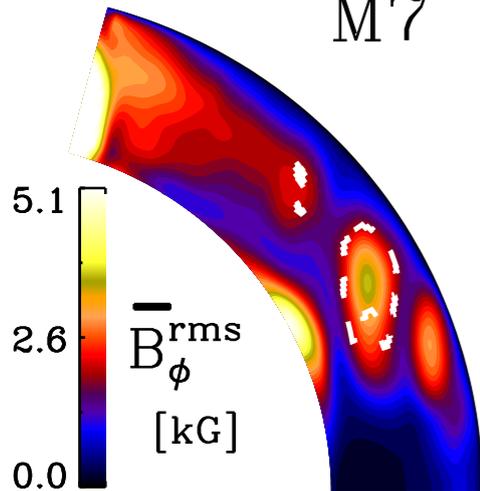
Found in observations and simulations

Berdyugina & Tuominen, 1998 & Lehtinen et al 2018

Viviani et al. 2018

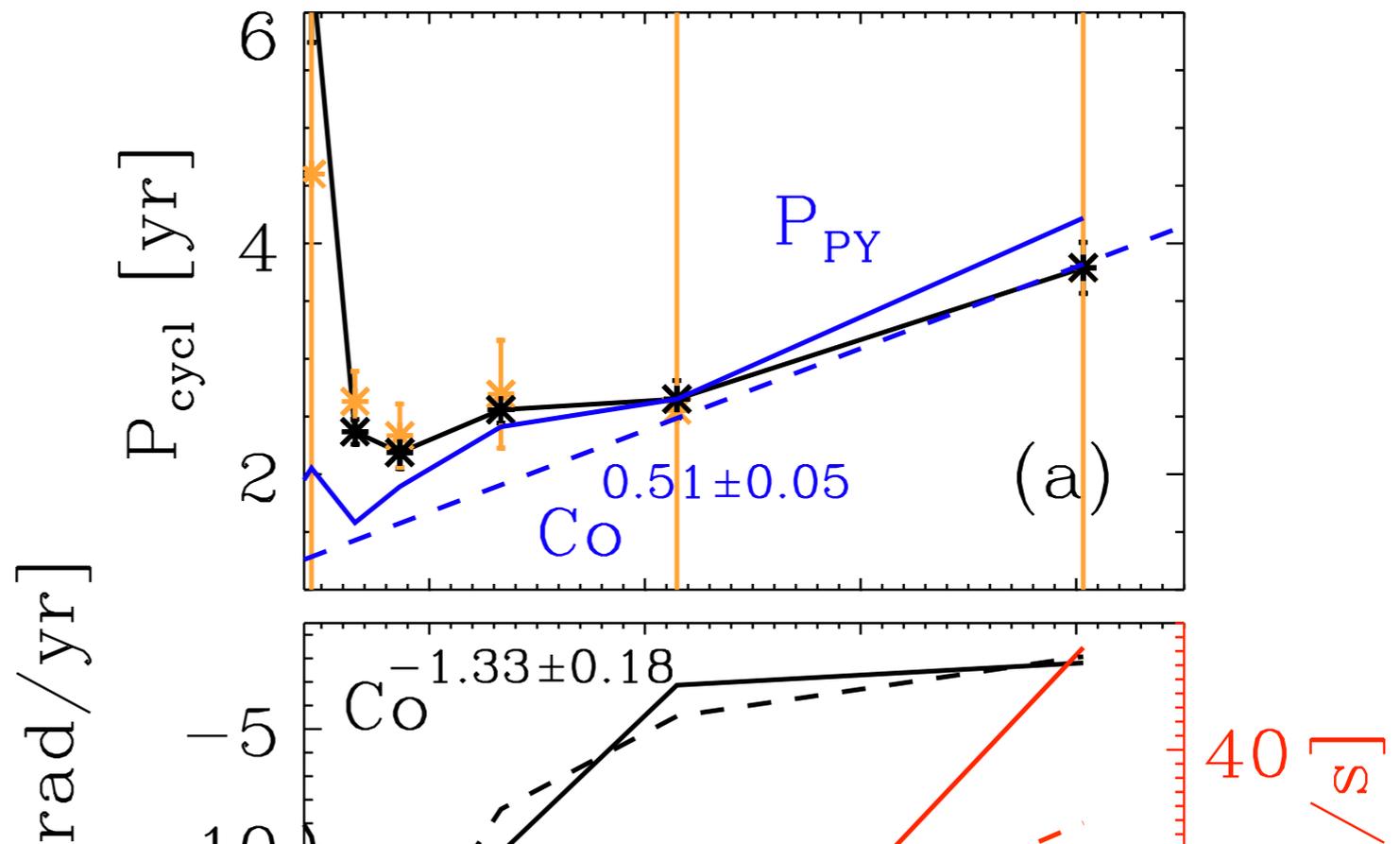
Dynamo cycles in simulations II

M7



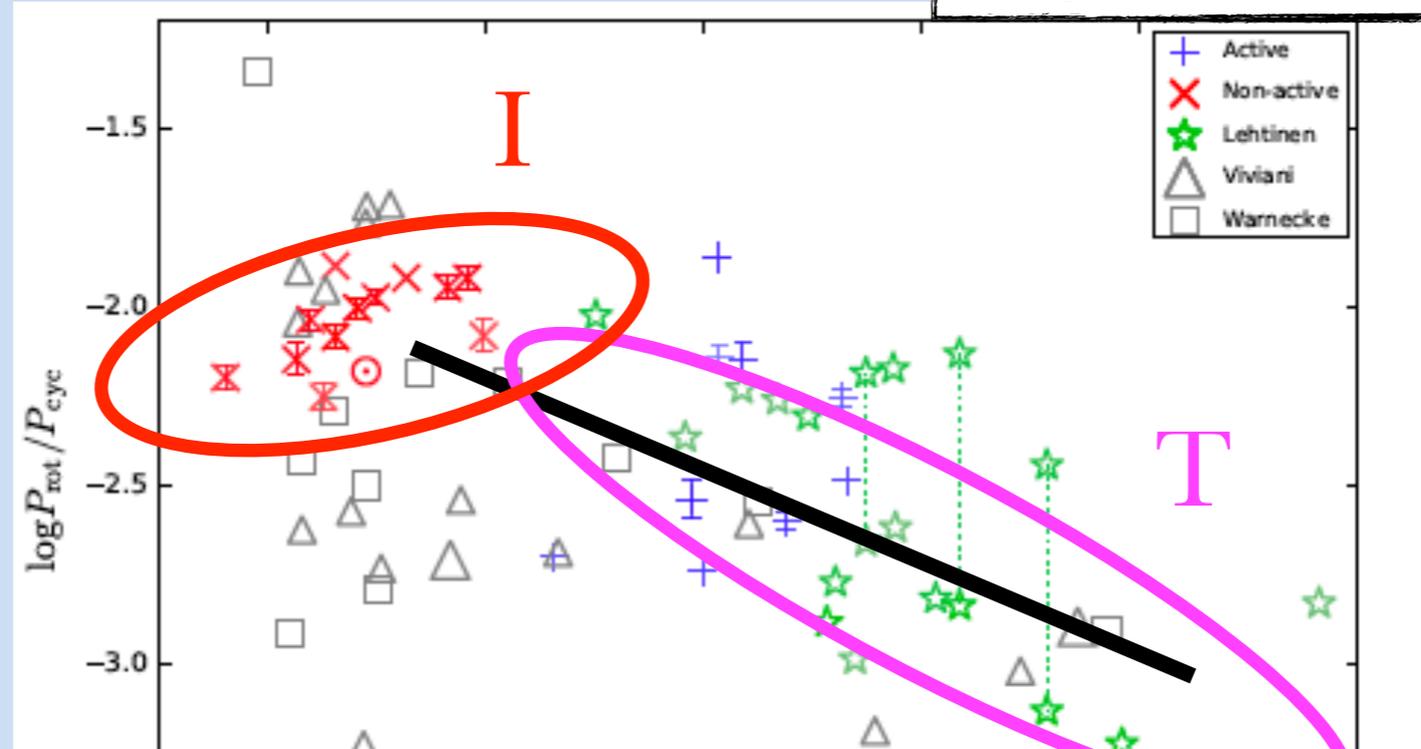
$$\omega_{\text{PY}} = \left| \frac{\alpha_{\phi\phi} k_{\theta}}{2} r \cos \theta \frac{\partial \Omega}{\partial r} \right|^{1/2}$$

Parker-Yoshimura



Scaling of cycle periods can be well described with P-Y dynamo wave.

Olsper et al. 2017



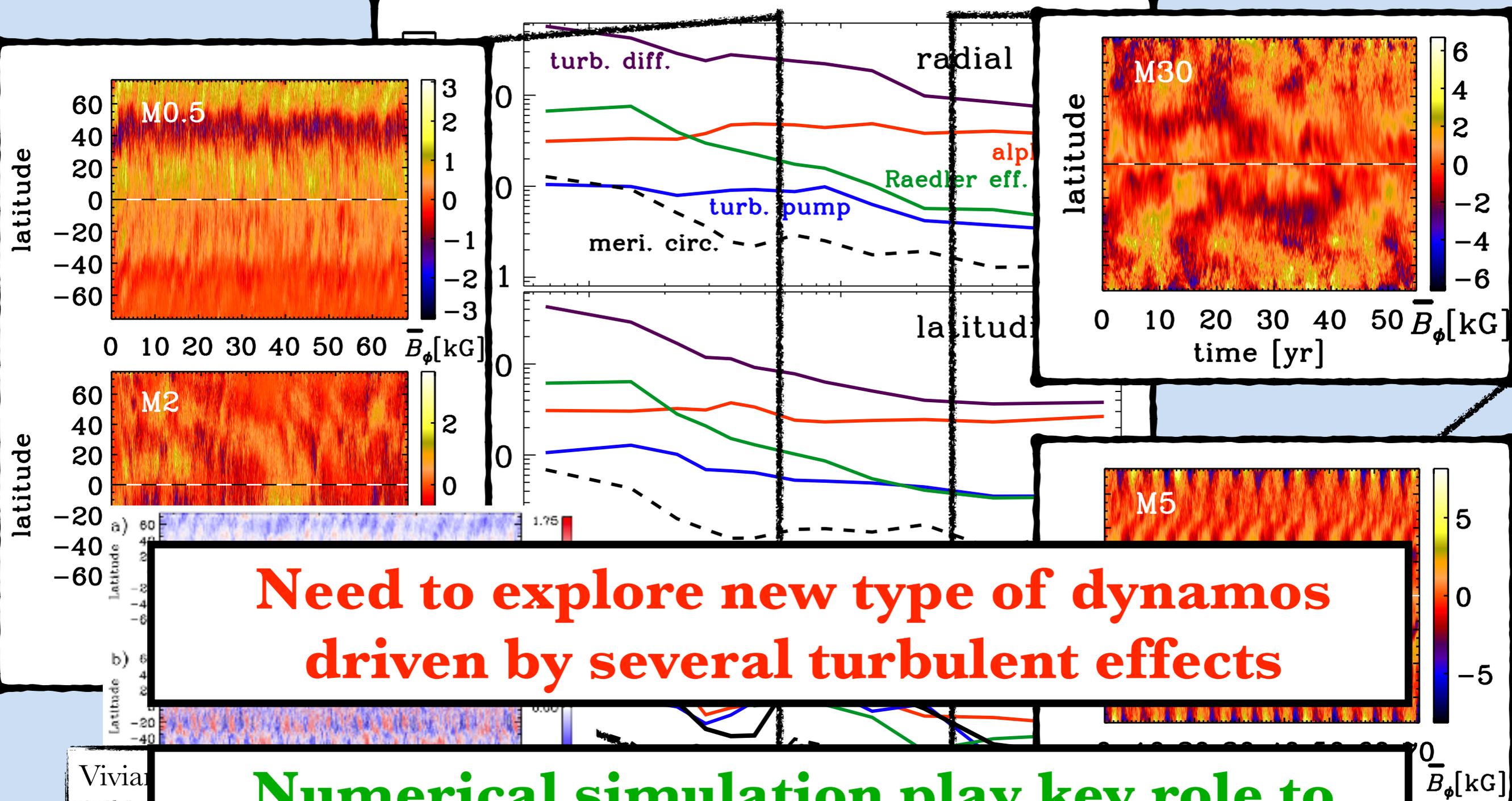
Results of simulations fit well with transitional branch

Transitional branch

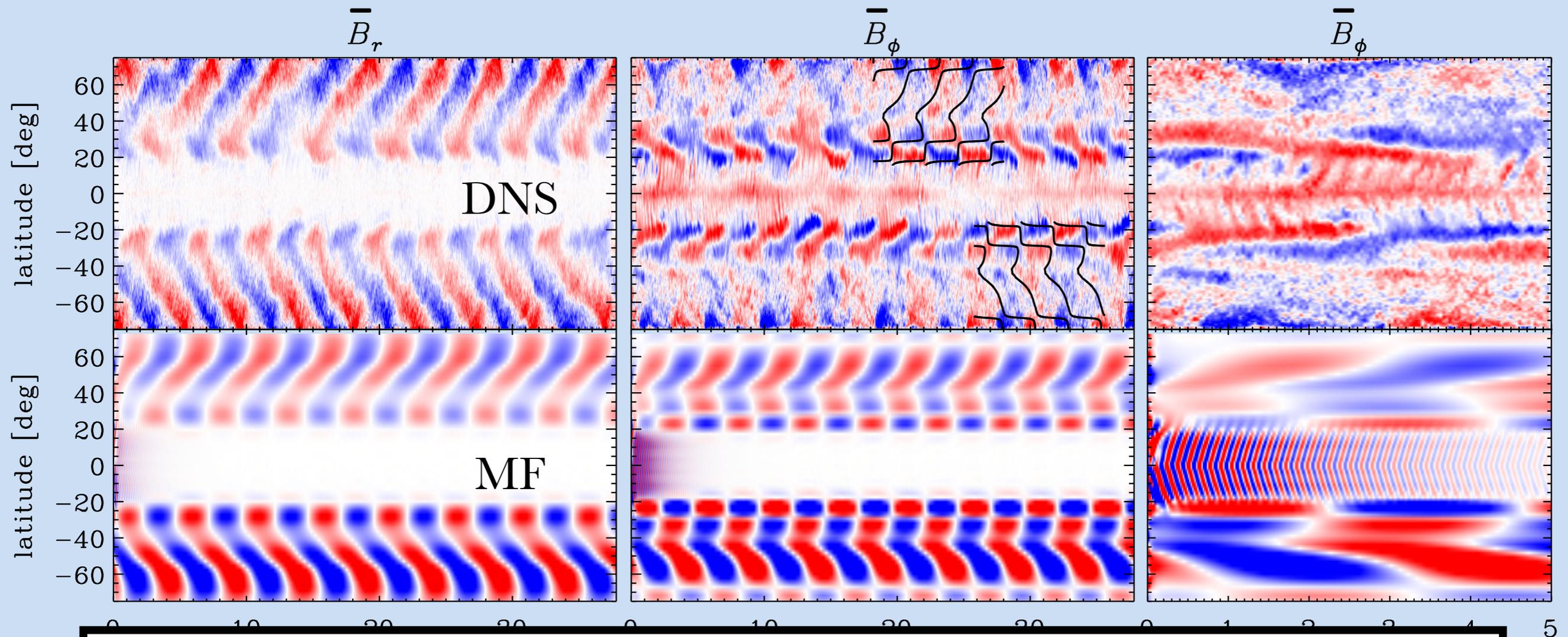
Distefano et al. 2017

Explanation for inactive branch still missing

Turbulent dynamo effects



Mean-field model vs. DNS



**Very good agreement of MF model with DNS:
(Period and Pattern)**

Full spectrum of effects needed

Simple dynamo models may not viable

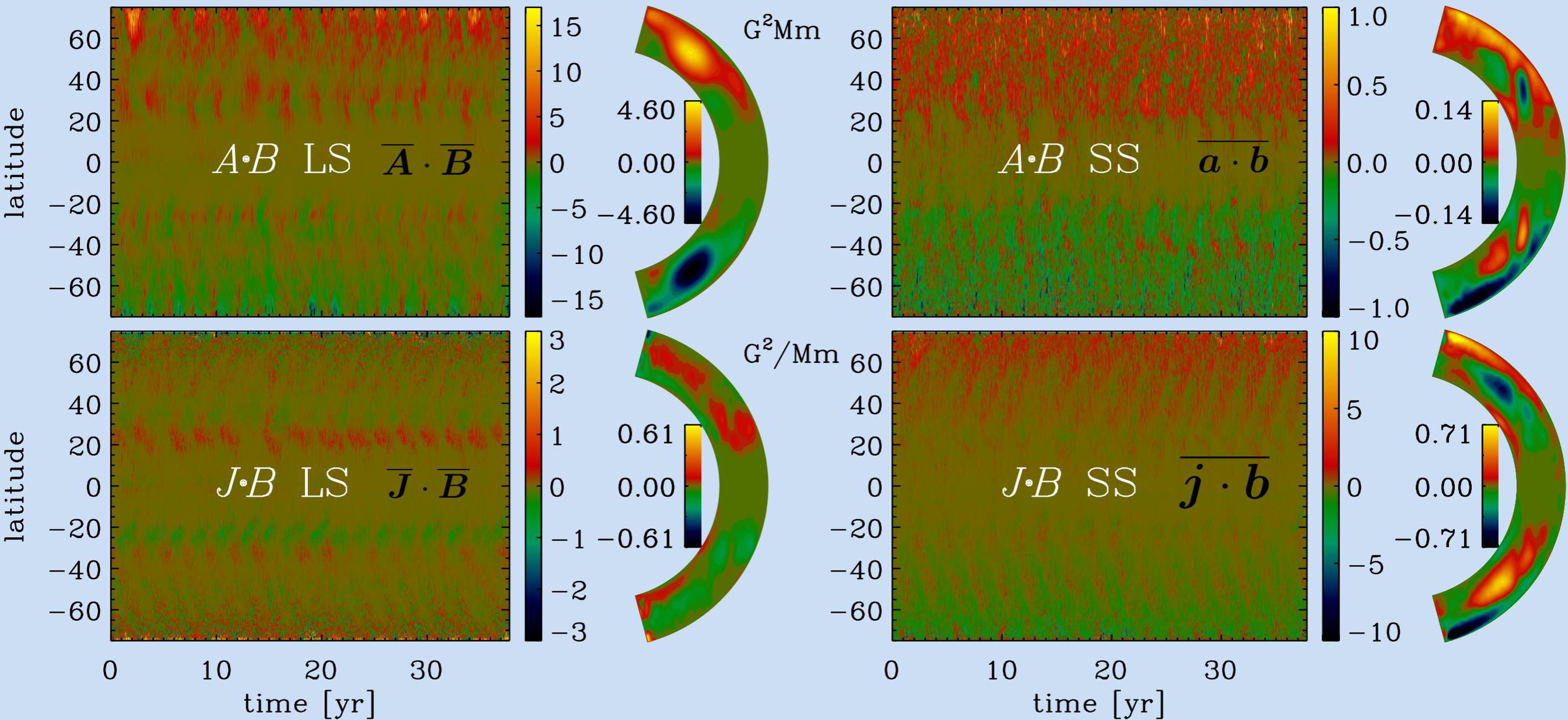
Warnecke et al. 2021
Submitted
arXiv:2105.07708

Magnetic helicity in convective dynamo simulation

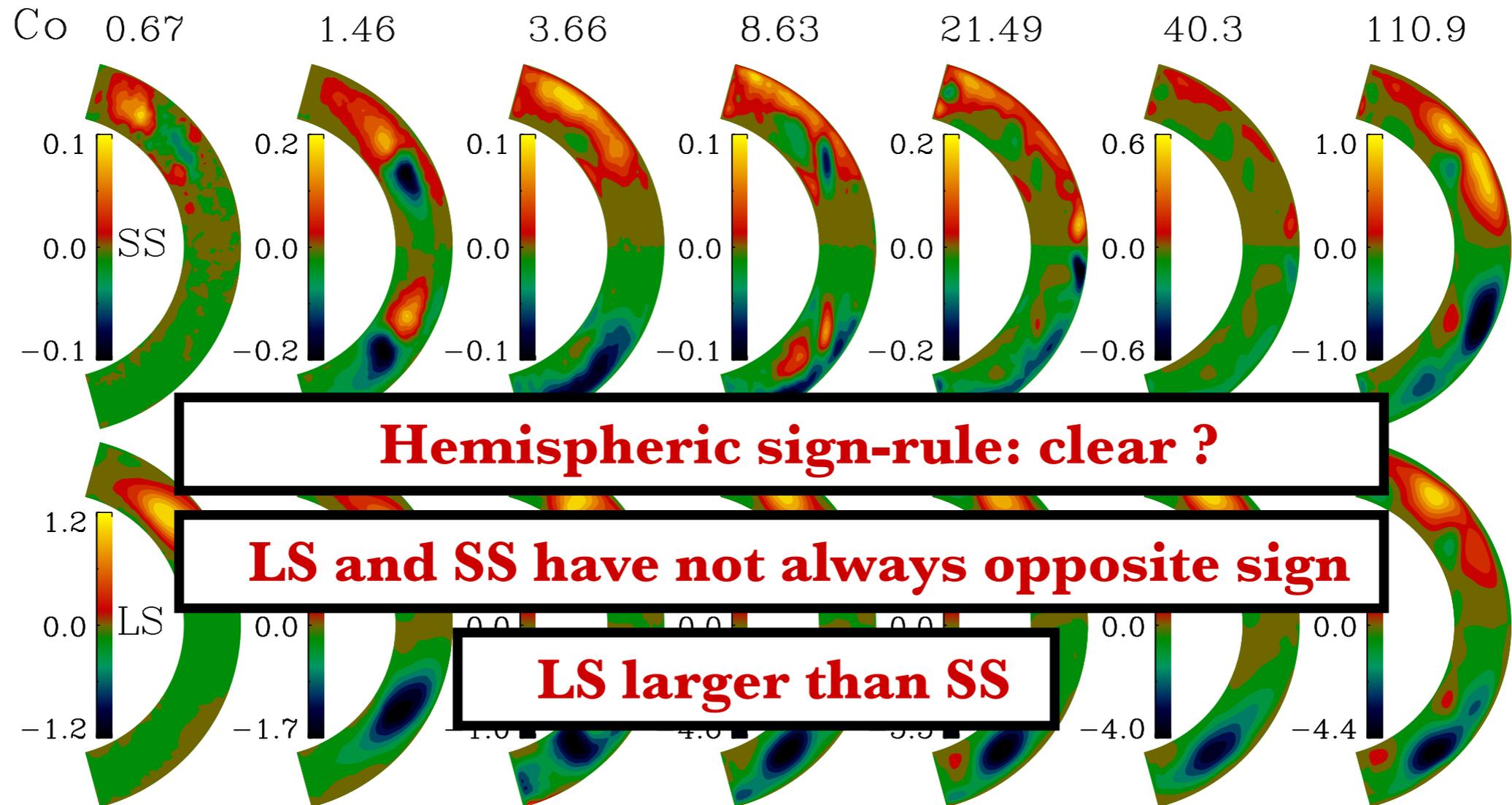
M5: $Co=8.63$

$r=0.98 R$

$r=0.98 R$



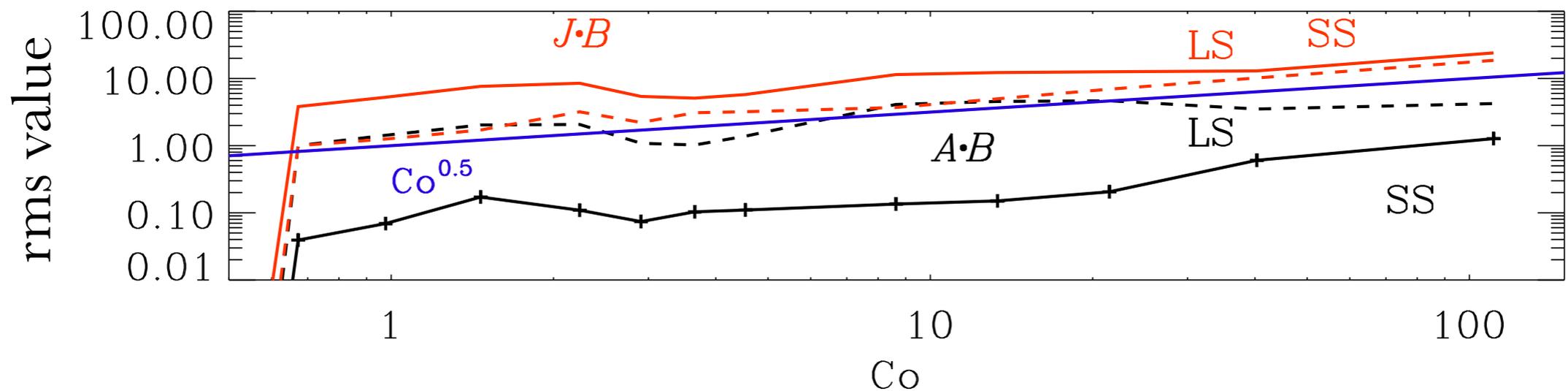
Magnetic helicity production with increasing rotation



Hemispheric sign-rule: clear ?

LS and SS have not always opposite sign

LS larger than SS



Magnetic helicity fluxes

$$\alpha = -\frac{1}{3} \tau_c \overline{\boldsymbol{\omega}' \cdot \mathbf{u}'} + \frac{1}{3} \frac{\tau_c}{\bar{\rho}} \overline{\mathbf{J}' \cdot \mathbf{B}'} = \alpha_K + \alpha_M,$$

Pouquet et al. 1976

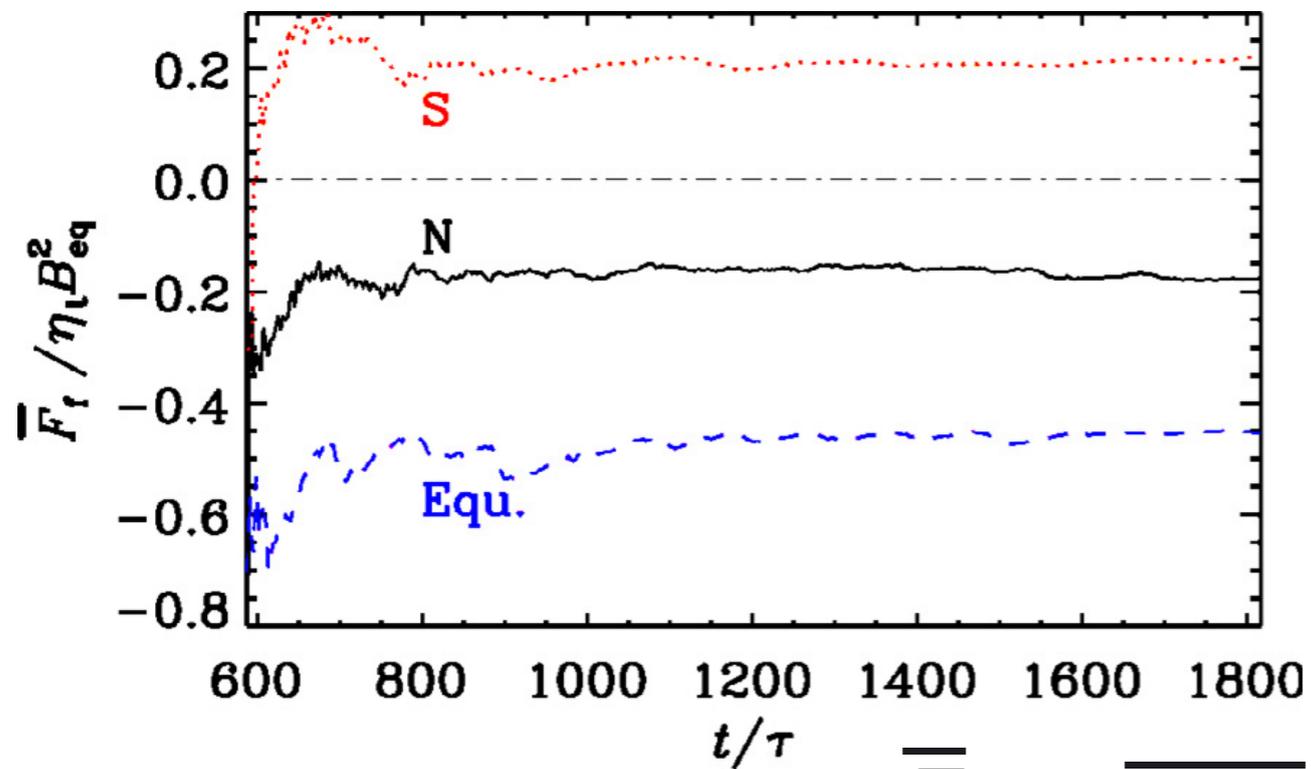
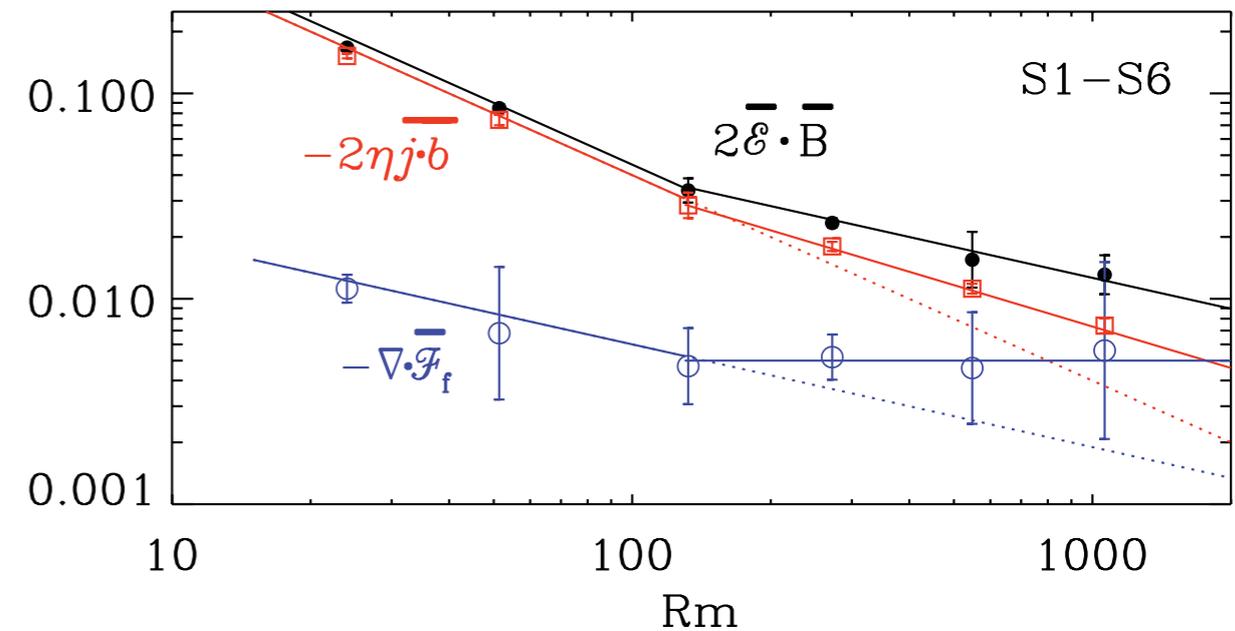
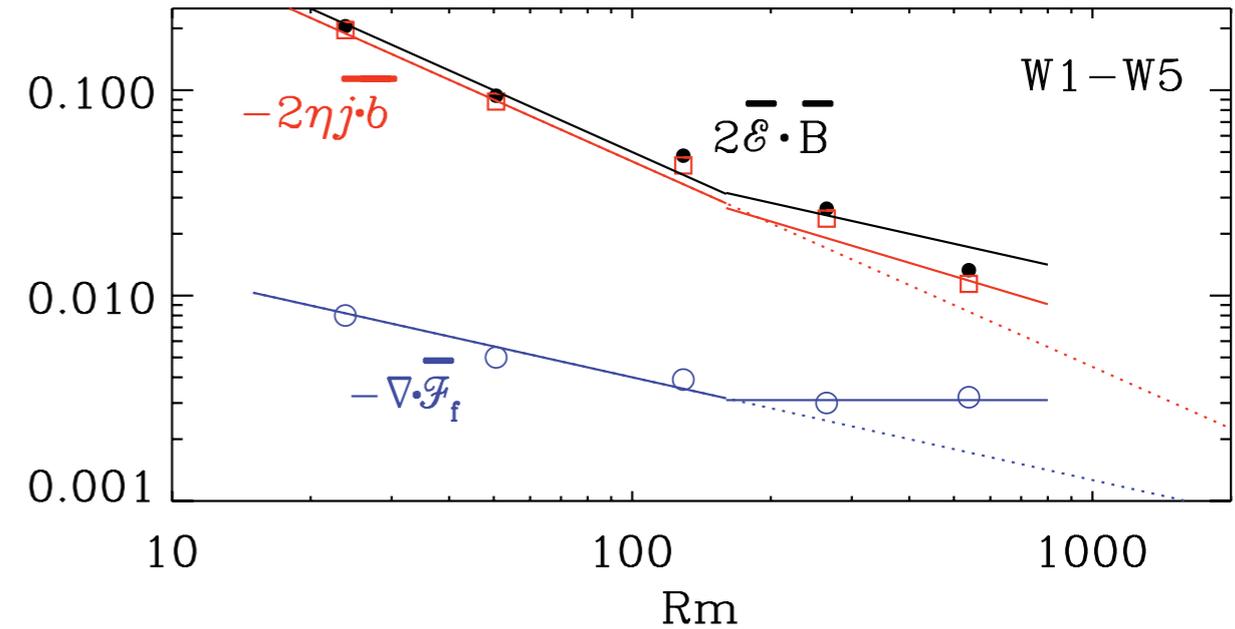
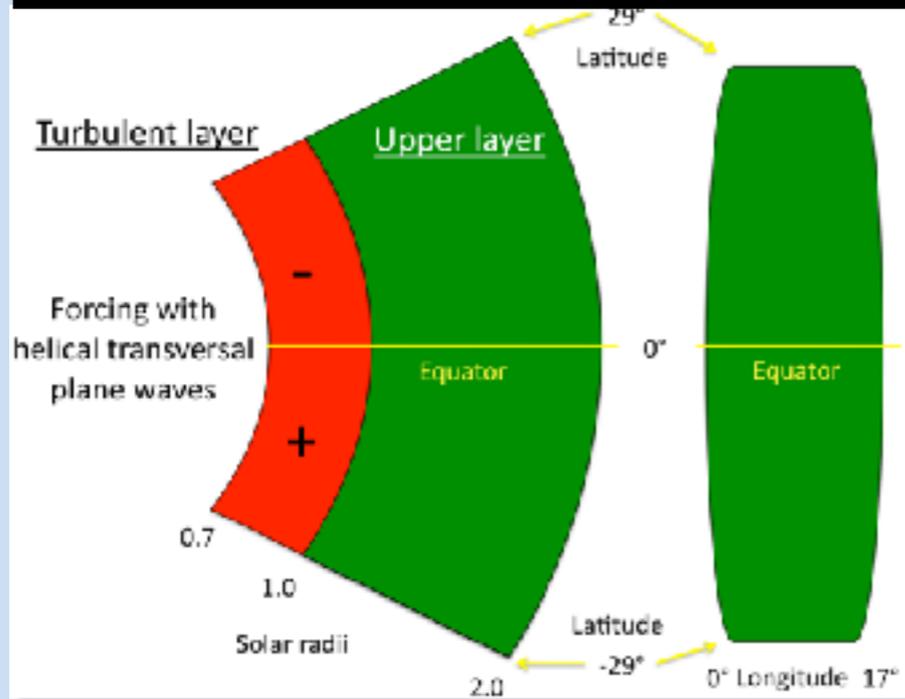
$$\alpha = \frac{\alpha_K + \text{Re}_M \left(\eta_t \bar{\mathbf{J}} \cdot \bar{\mathbf{B}} - \frac{1}{2} \nabla \cdot \bar{\mathcal{F}}_h^f \right) / B_{\text{eq}}^2}{1 + \text{Re}_M \bar{\mathbf{B}}^2 / B_{\text{eq}}^2}$$

Kleeorin & Ruzmaikin 1982

Vainshtein & Cattaneo 1992

**Based on assumption, which might be not true
in complex systems**

Magnetic helicity fluxes from simulations

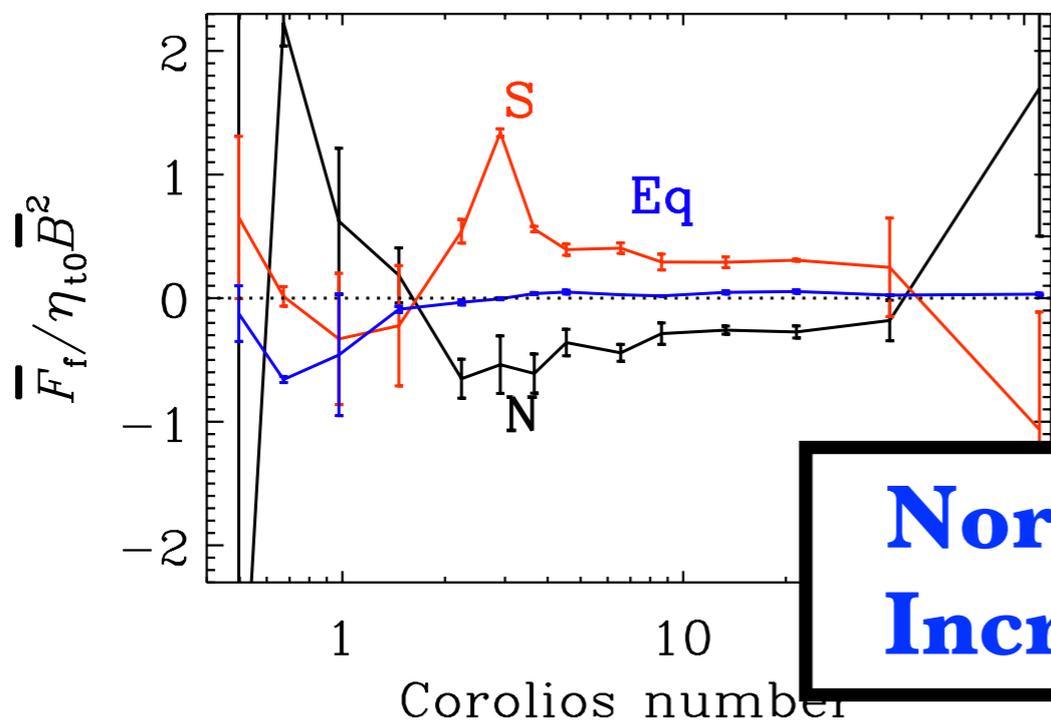
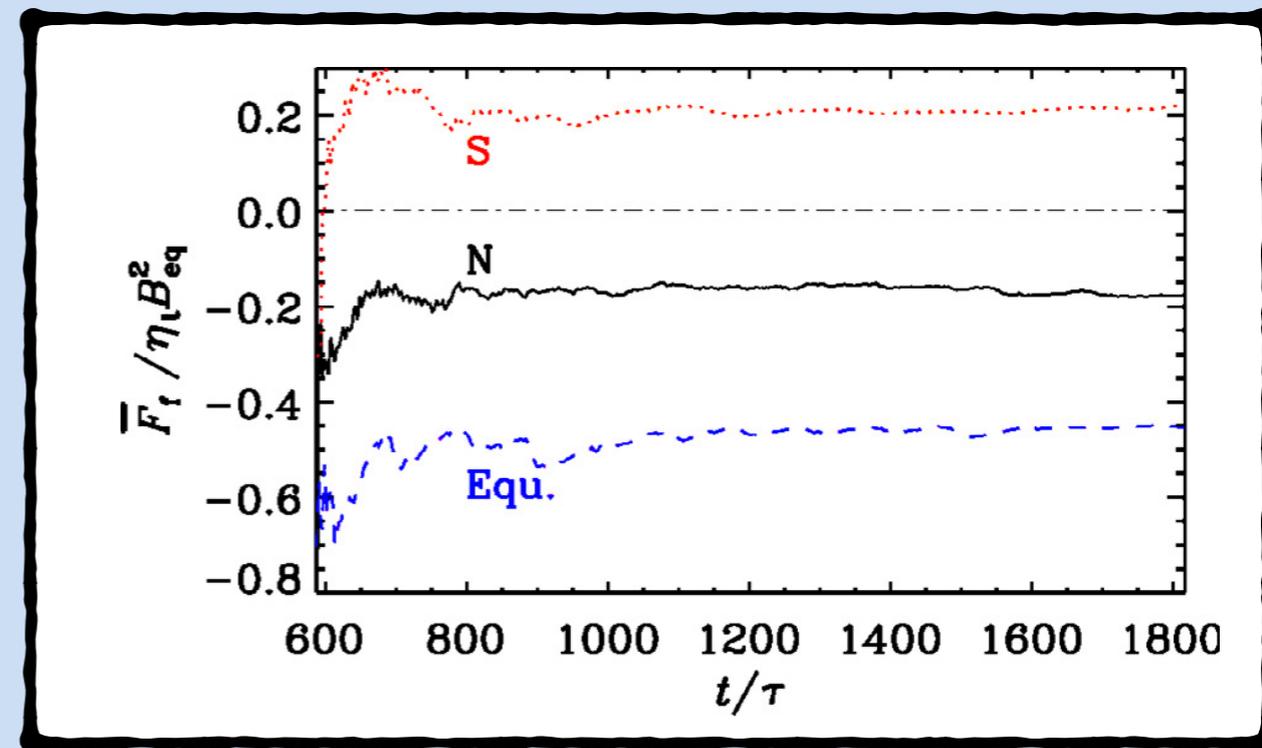
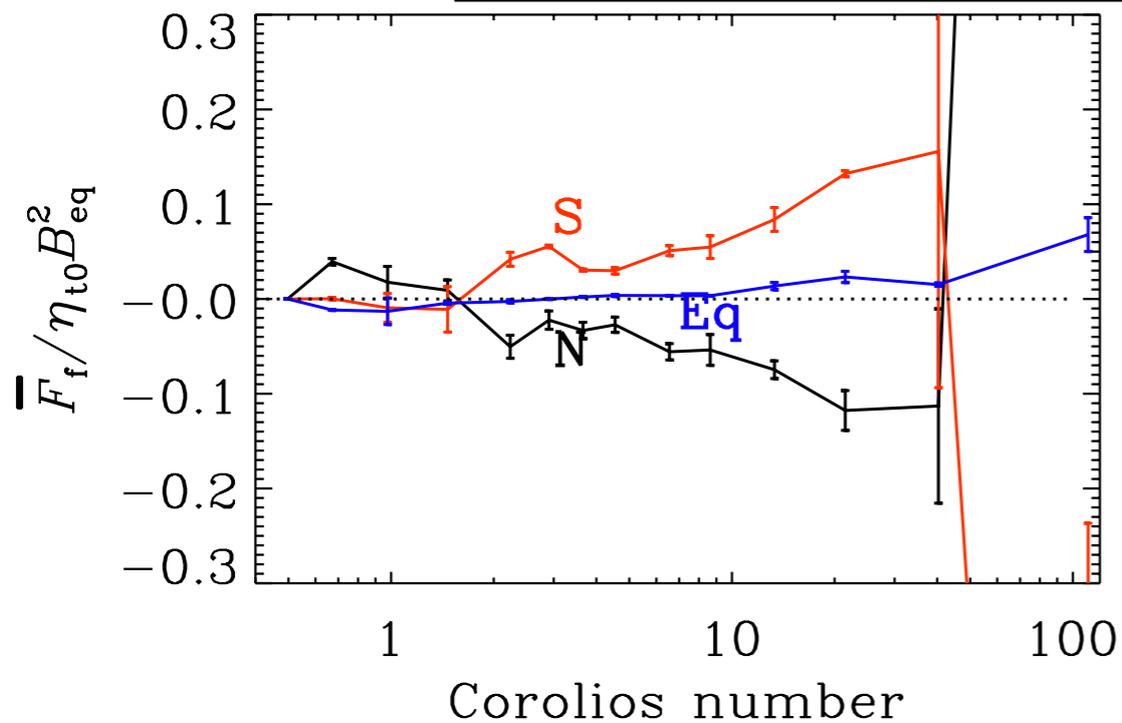


Warnecke et al. 2011

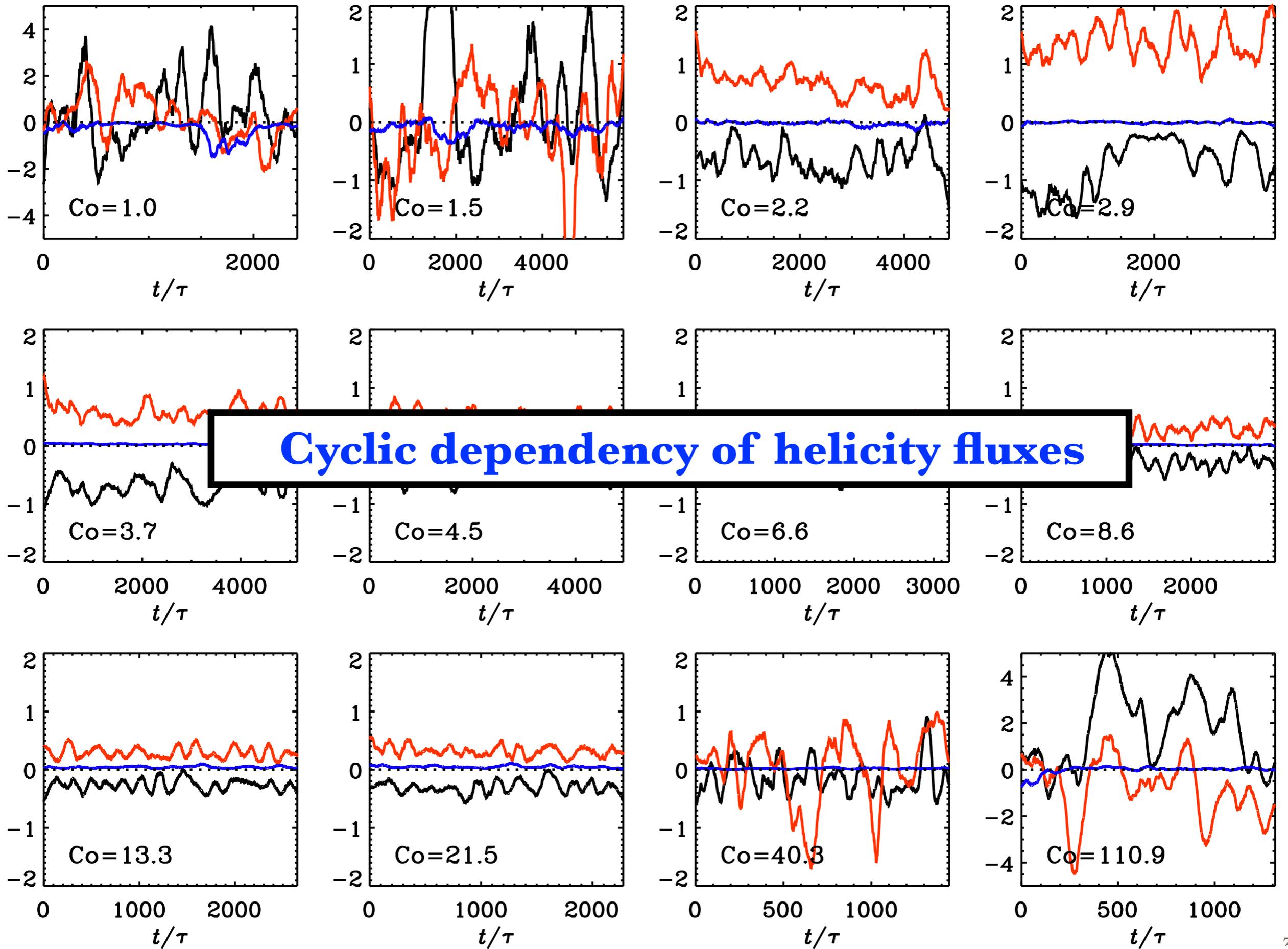
$$\bar{\mathbf{F}}_f = \overline{\mathbf{e} \times \mathbf{a}}$$

Del Sordo et al. 2013

Magnetic helicity fluxes from convection simulations



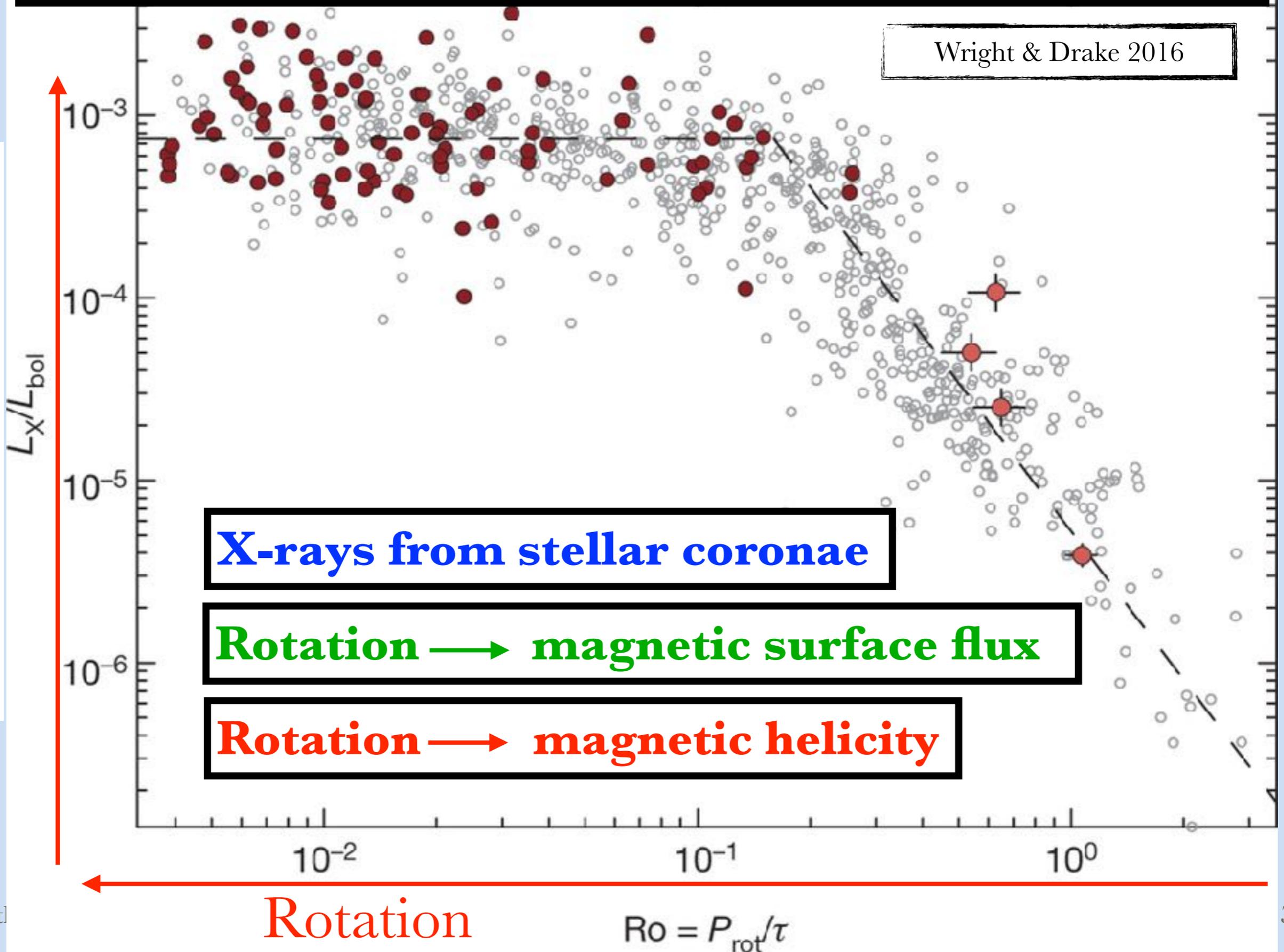
**Norm. helicity fluxes stay constant
Increase due to mag. field increase**



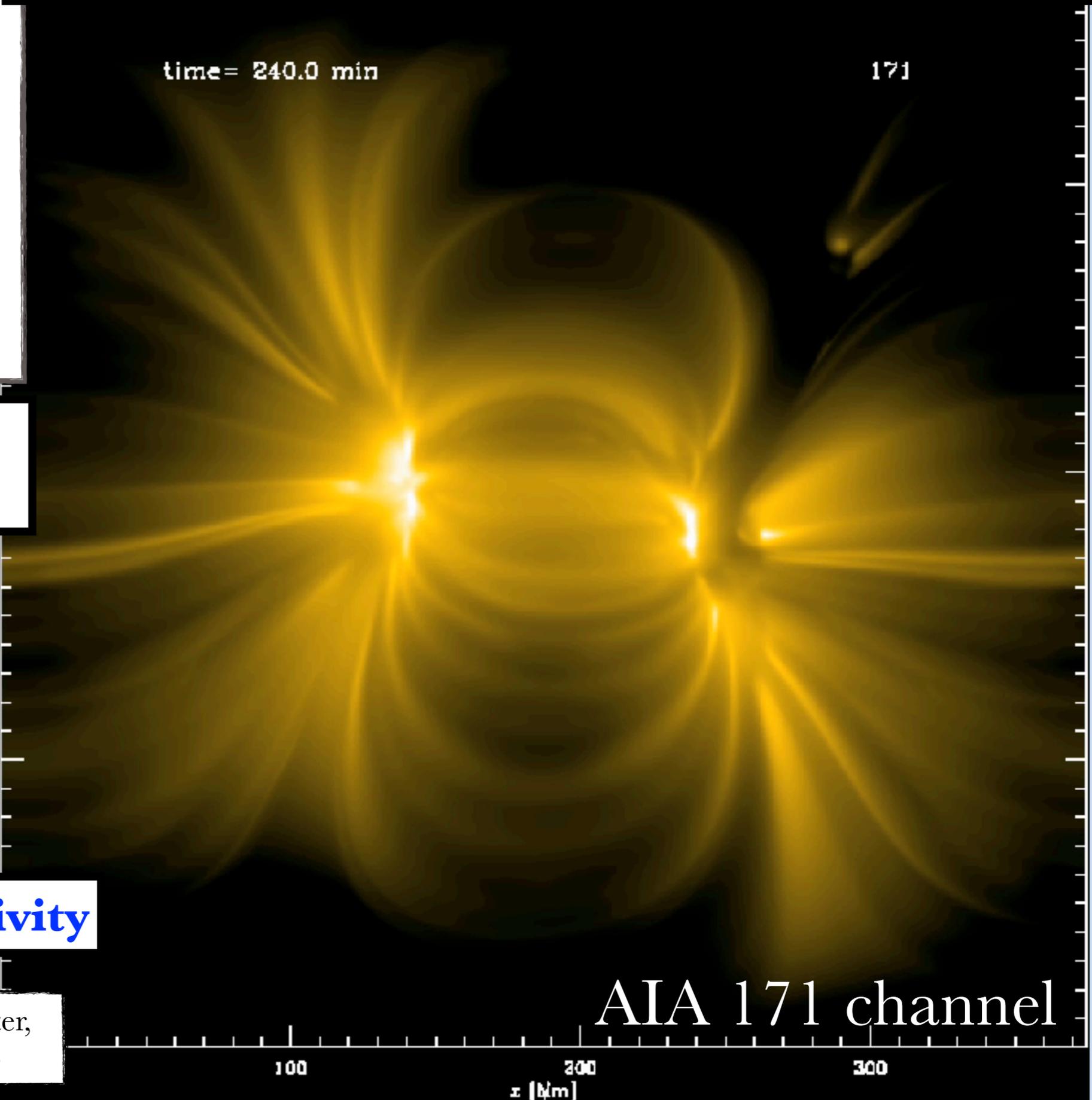
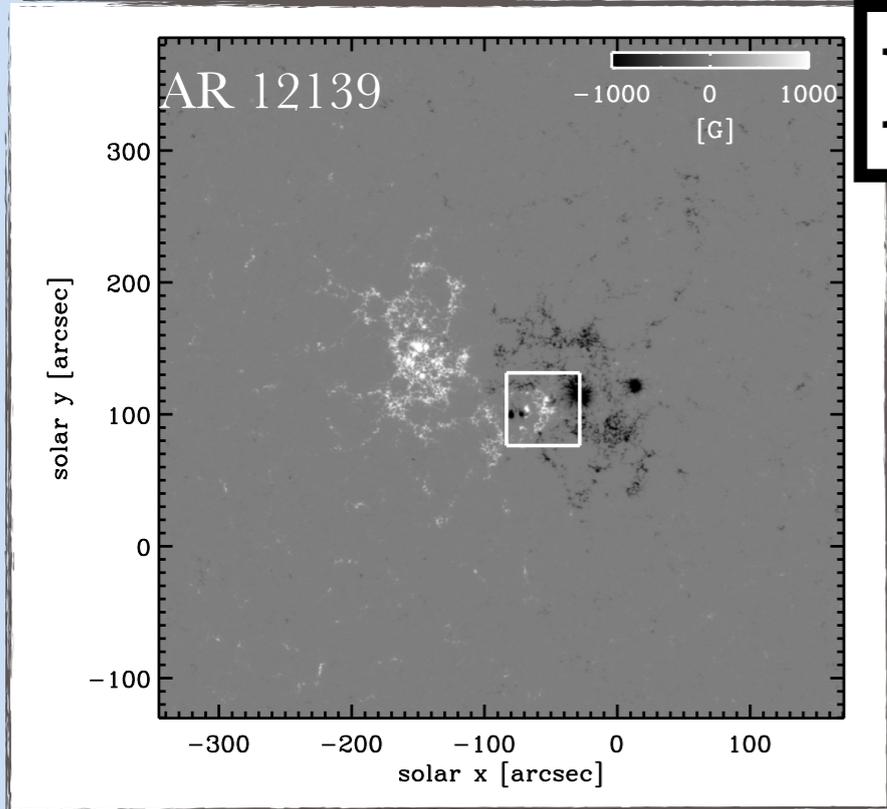
Rotation-Activity Relation

Wright & Drake 2016

X-ray Luminosity



Data-driven coronal simulation



**Compressible
(resistive) MHD**

**Prescribed
photospheric
motions
mimicking solar
granulation**

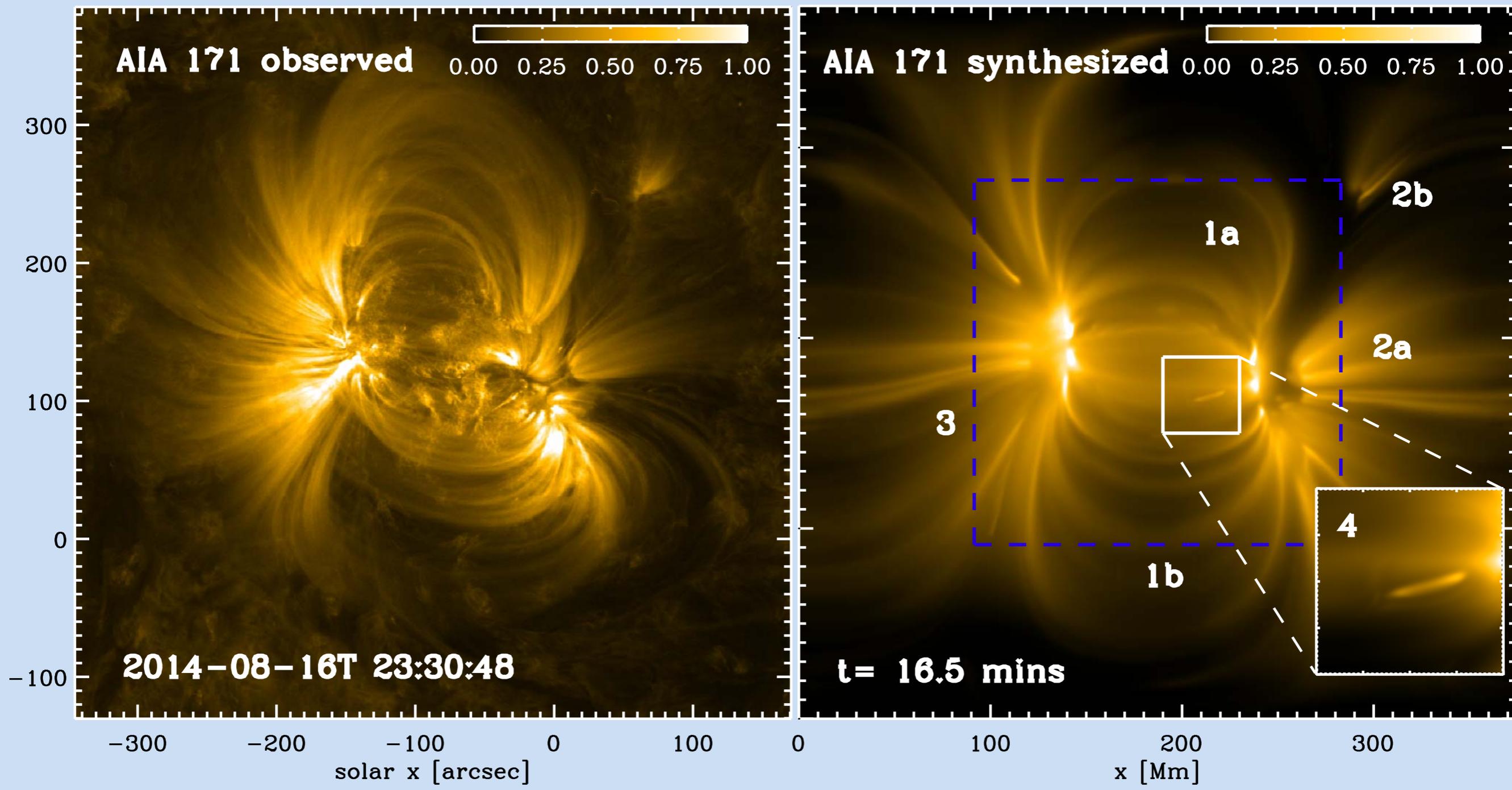
Spitzer heat conductivity



Bingert & Peter,
2011, 2013

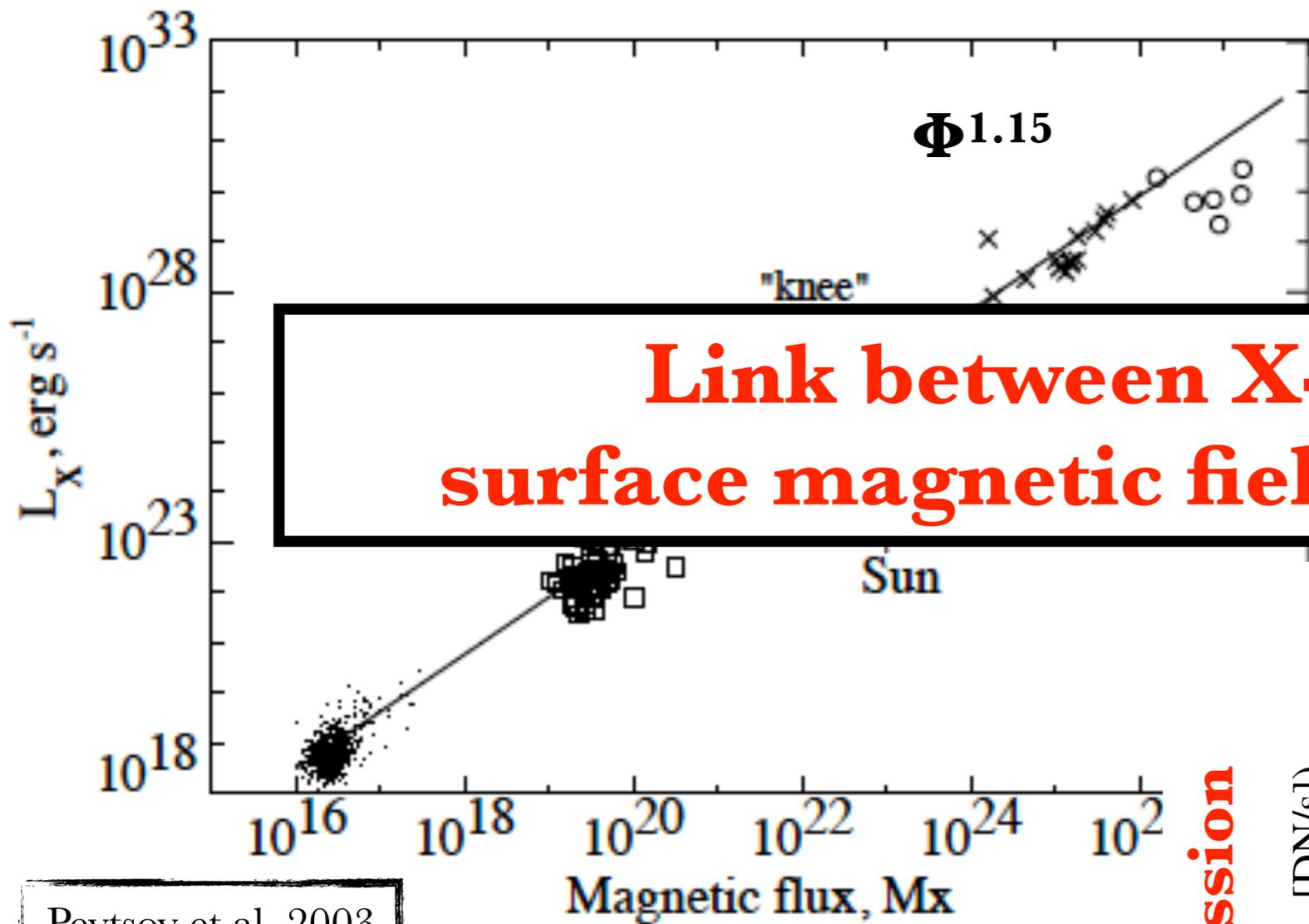
AIA 171 channel

Comparison with observations



Can reproduce many observed features !

Scaling of X-rays with magnetic flux



Link between X-rays and surface magnetic field distribution

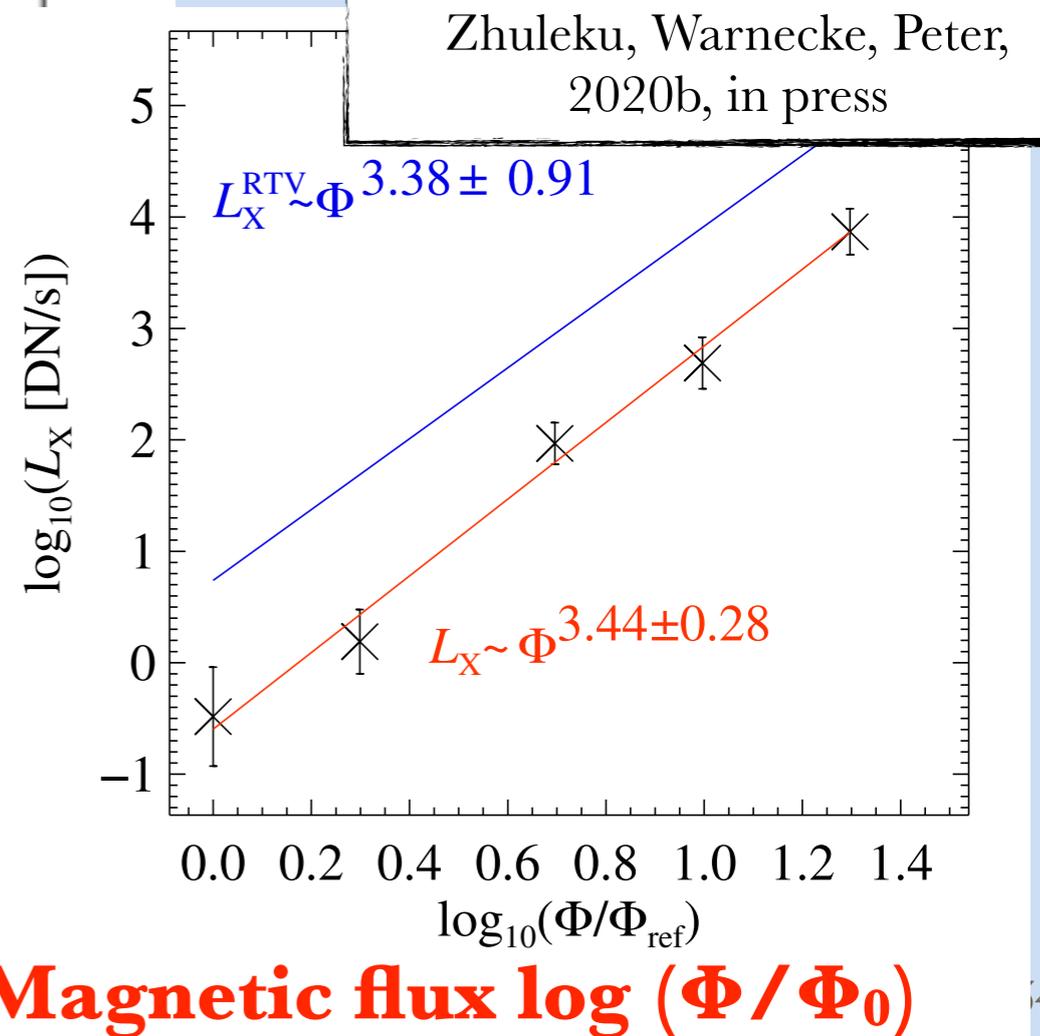
Pevtsov et al. 2003

See also Vidotto et al. 2014

Explained by analytical model

Zhuleku, Warnecke, Peter, 2020a

X-ray emission

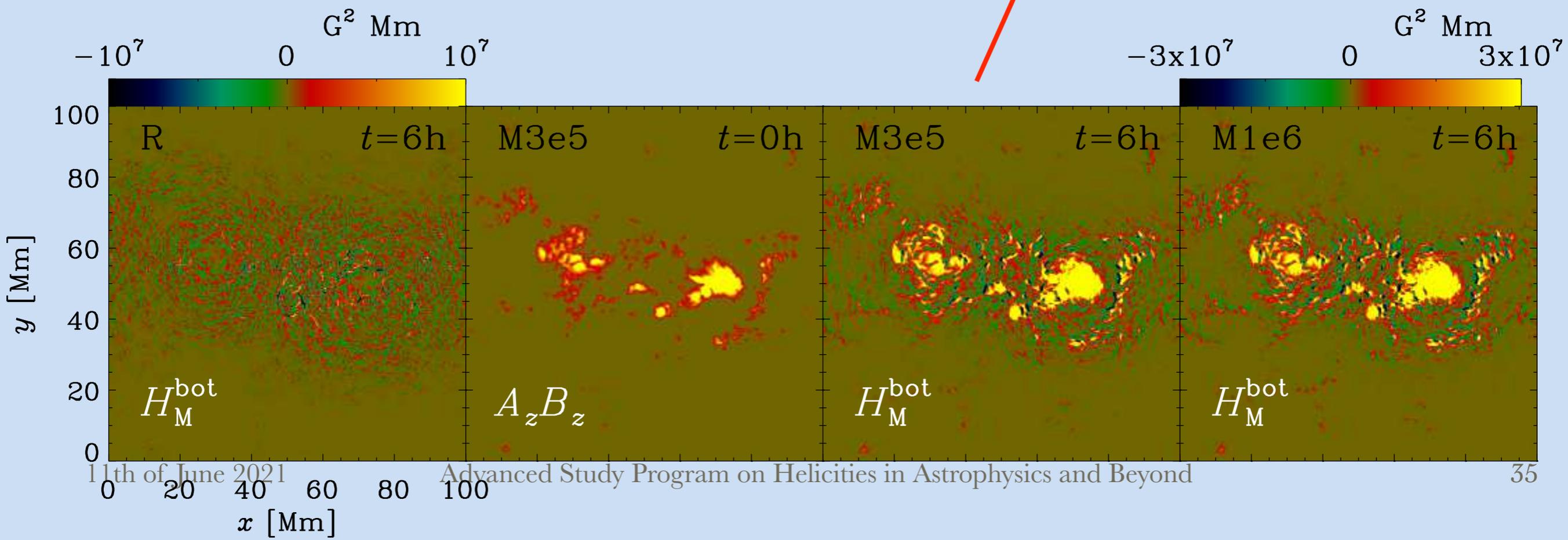
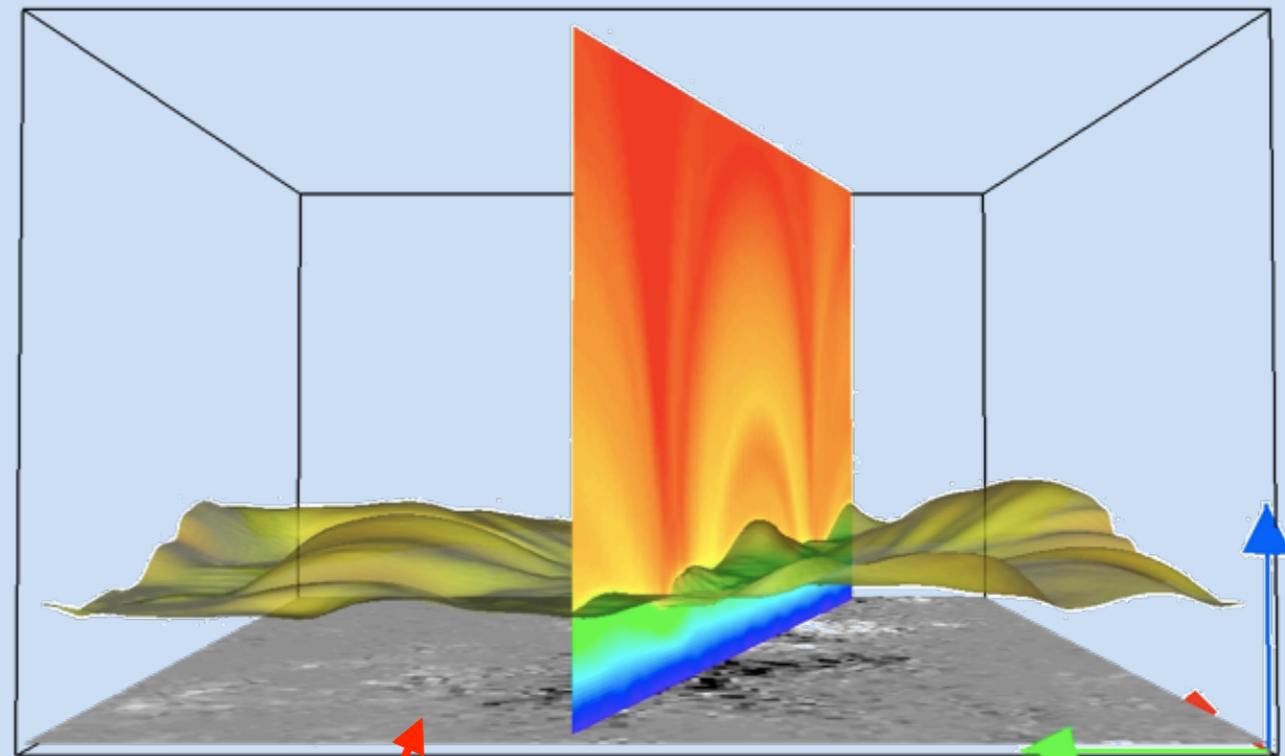


Magnetic helicity injection

Vertical photospheric field: A_x, A_y

$$A_z = \frac{H_M^{\text{in}}}{\langle B_z^2 \rangle_{xy}} B_z,$$

$$H_M^{\text{bot}} = \langle \mathbf{A} \cdot \mathbf{B} \rangle_{xy} (z = 0).$$



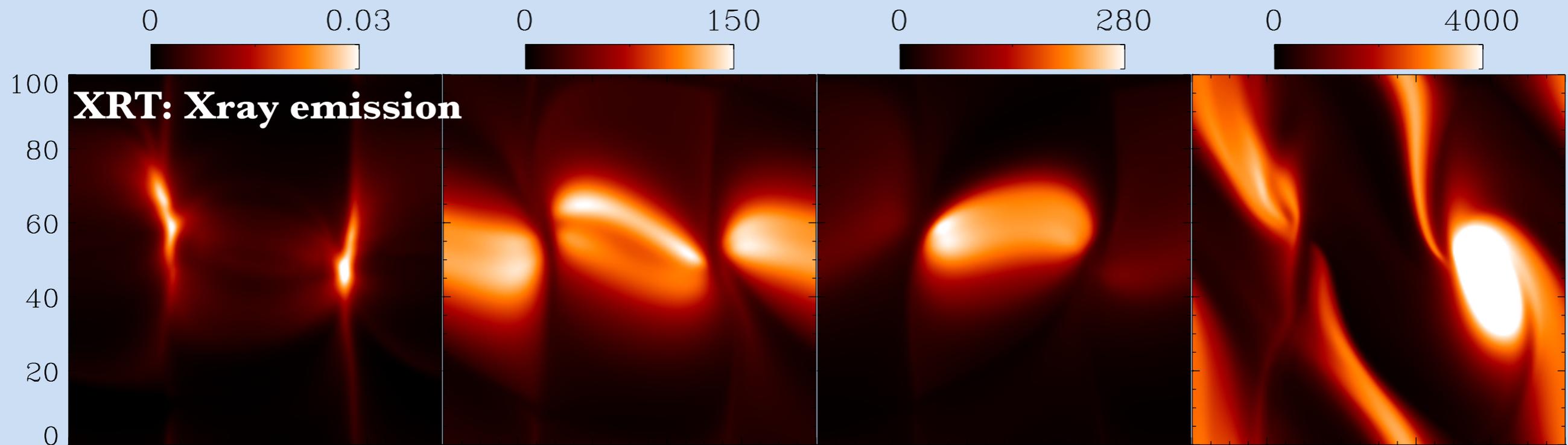
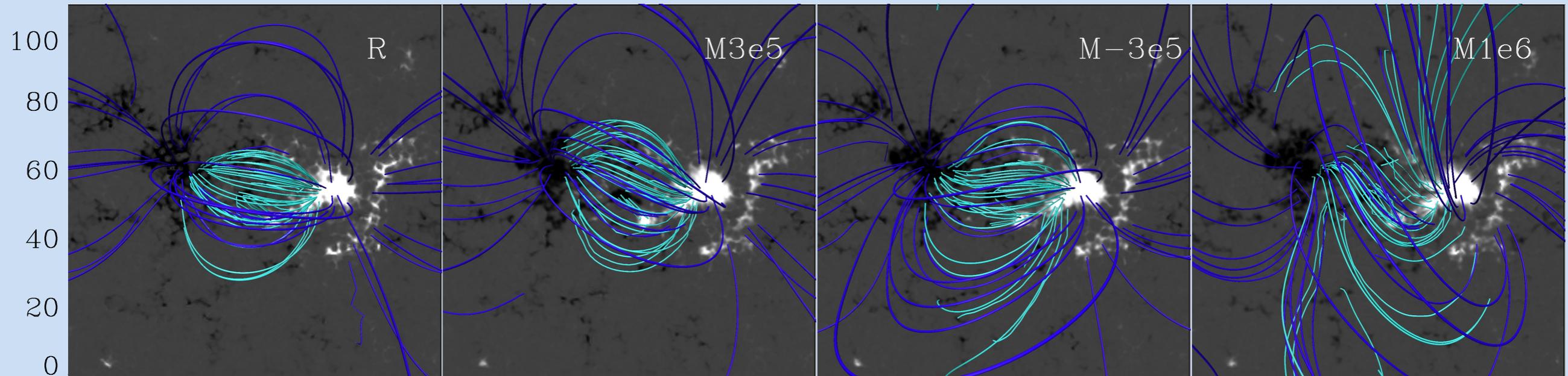
Magnetic helicity + X-ray production

$H_m=0$

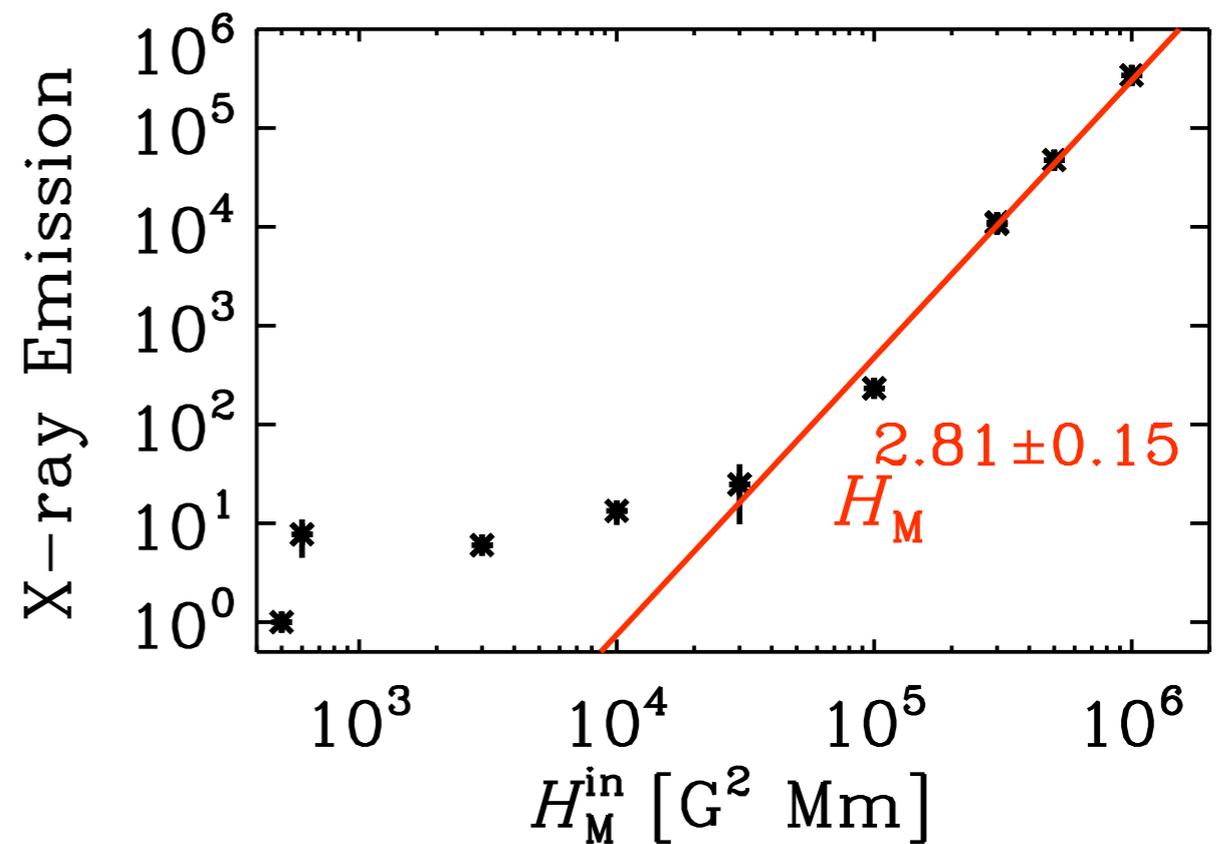
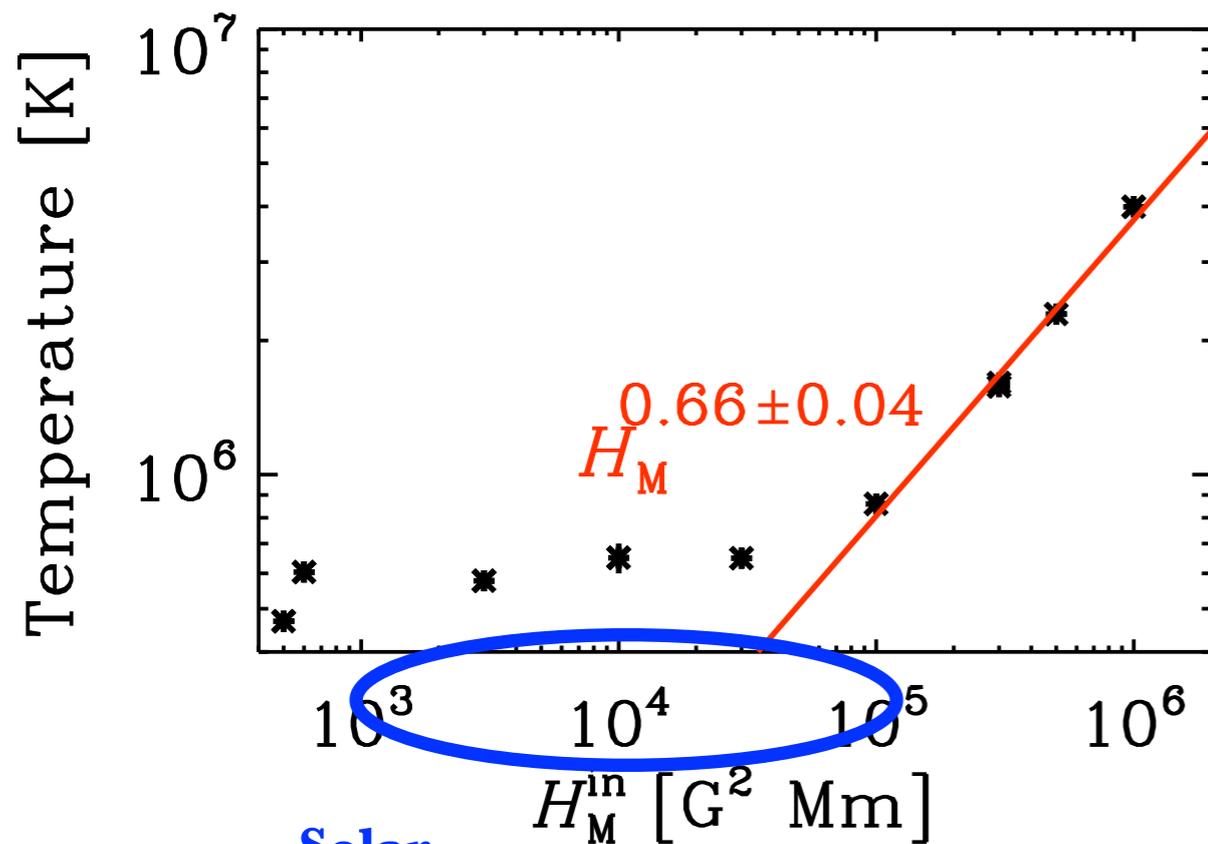
$H_m=3 \times 10^5 \text{ G}^2 \text{ Mm}$

$H_m=-3 \times 10^5 \text{ G}^2 \text{ Mm}$

$H_m=1 \times 10^6 \text{ G}^2 \text{ Mm}$



Magnetic helicity enhances heating



Magnetic helicity density

Increase of X-rays with magnetic helicity is consistent with observation, if H_m increases less linearly with rotation.

Increase of X-rays with magnetic helicity provide a significant contribution to activity-rotation-relation

Conclusions

Test-field method one way to measure dynamo effects

Trace of alpha saturates for high rotation, but non-isotropic

Explanation for inactive branch still missing

Large variety of turbulent dynamo effects

Simple dynamo models may not viable

LS and SS mag. helicity no opposite sign

All helicities follow $\sim Co^{0.5}$

Norm. helicity fluxes independent of rotation

Injecting mag. helicity causes higher Xray fluxes

Significant contribution to Xray rotation relation

Helical magnetic field important role in understanding the Rotation-Activity-Relation of stars